

SUMMARY OF LECTURE #4:

• BGK EQUATION = $\partial_t f + v \cdot \nabla_x f = Q(f)$
 $= \frac{1}{2} (M[\rho, u, T] - f)$

• DISCRETE VELOCITIES:

$v_k \in \text{GRID } V$ (N DISCRETE VELOCITIES)

• $f(t, x, v_k) \approx F_k(t, x)$

$F(t, x) = (F_k(t, x))$ IS A DISCRETE VELOCITY DISTRIBUTION FUNCTION

• $\begin{pmatrix} \rho \\ p_y \\ E \end{pmatrix} = \sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} F_k(t, x) \omega_k$

• DISCRETE VELOCITY APPROXIMATION OF BGK:

|| $\partial_t F_k + v_k \cdot \nabla_x F_k = Q_k(F)$
 $= \frac{1}{2} (\underbrace{M[\rho, u, T]}_{\text{TO BE DEFINED}} - F_k)$

HOW TO DEFINE THE APPROXIMATION $(M_h(\rho, u, T))$
OF $M(\rho, u, T)$?

↳ CONSTRAINT: PRESERVE THE CONSERVATION PROPERTIES:

$$\int_{\mathbb{R}^3} \begin{pmatrix} 1 \\ v \\ \frac{1}{2} \|v\|^2 \end{pmatrix} Q(f) dv = 0$$

THAT ARE EQUIVALENT TO :

$$\int_{\mathbb{R}^3} \begin{pmatrix} 1 \\ v \\ \frac{1}{2} \|v\|^2 \end{pmatrix} M(\rho, u, T)(v) dv = \begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}$$

(WHICH IS TRUE)

↳ AT THE DISCRETE LEVEL, WE WANT:

$$\left[\sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} M_k(\rho, u, T) w_k = \begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix} \right.$$

("CONSERVATIVE APPROXIMATION")

BUT THE NATURAL APPROXIMATION:

$$M_k[\rho, u, \tau] = M[\rho, u, \tau](\omega_k) \quad \text{IS } \underline{\text{NOT CONSERVATIVE}}$$

$$\begin{aligned} \text{INSTEAD: } & \sum_{k=1}^N M_k[\rho, u, \tau] \omega_k \\ &= \sum_{k=1}^N M[\rho, u, \tau](\omega_k) \omega_k \end{aligned}$$

WHICH IS ONLY AN APPROXIMATION OF

$$\int_{\mathbb{R}^3} M[\rho, u, \tau](\omega) d\omega = \rho$$

$$\text{THEREFORE: } \sum_{k=1}^N M[\rho, u, \tau](\omega_k) \omega_k \neq \rho$$

CONSERVATIVE APPROXIMATION :

DEFINE PARAMETERS $\tilde{\rho}, \tilde{u}, \tilde{T}$ SUCH THAT

$$\begin{aligned} M_k[\rho, u, T] &= M[\tilde{\rho}, \tilde{u}, \tilde{T}](\psi_k) \\ &= \frac{\tilde{\rho}}{(2\pi\tilde{T})^{3/2}} \exp\left(-\frac{\|\psi_k - \tilde{u}\|^2}{2\tilde{T}}\right) \end{aligned}$$

SATISFIES :

$$\sum_{k=1}^N \begin{pmatrix} 1 \\ \psi_k \\ \frac{1}{2} \|\psi_k\|^2 \end{pmatrix} M_k[\rho, u, T] \omega_k = \begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}$$

CONSERVATION PROPERTIES
ARE ENFORCED

L) MATHEMATICAL ANALYSIS : SEE
L. MIEUSSONS (MBAS, VOL. 8 (10), 2000)

$$\sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} M_k(\rho, u, T) w_k = \begin{pmatrix} \rho \\ p_u \\ E \end{pmatrix}$$

$$= \sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} M[\tilde{\rho}, \tilde{u}, \tilde{T}] (v_k) w_k$$

"THERE EXISTS
A DISCRETE DISTRIBUTION
FUNCTION G . S.T
 $\sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} G_k w_k = \begin{pmatrix} \rho \\ p_u \\ E \end{pmatrix}$
AND $G_k > 0$ "

THIS IS A NON-LINEAR SYSTEM OF 5 EQUATIONS
AND 5 UNKNOWNNS ($\tilde{\rho}, \tilde{u}, \tilde{T}$).

↳ THERE EXISTS A UNIQUE SOLUTION
IF $(\tilde{\rho}, \tilde{u}, \tilde{T})$ SATISFY SOME CONSTRAINTS
AND \forall

↳ HOW TO COMPUTE THIS SOLUTION?

→ NO EXPLICIT SOLUTION

→ NUMERICAL APPROXIMATION
BY A NEWTON ALGORITHM

NON LINEAR SYSTEM: $\tilde{\rho}, \tilde{u}, \tilde{T}$ s.t

$$\sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} M[\tilde{\rho}, \tilde{u}, \tilde{T}](v_k) \omega_k = \begin{pmatrix} \rho \\ pu \\ E \end{pmatrix}$$

LET $\vec{\alpha} = \begin{pmatrix} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{pmatrix}$ s.t. $M[\tilde{\rho}, \tilde{u}, \tilde{T}](v) = \exp(\vec{\alpha} \cdot \vec{m}(v))$

$\vec{m} = \begin{pmatrix} 1 \\ v \\ \frac{1}{2} \|v\|^2 \end{pmatrix}$

$$M[\tilde{\rho}, \tilde{u}, \tilde{T}](v) = \frac{\tilde{\rho}}{(2\pi\tilde{T})^{3/2}} \exp\left(-\frac{\|v-u\|^2}{2\tilde{T}}\right)$$

$$\Leftrightarrow \alpha_1 = \log \frac{\tilde{\rho}}{(2\pi\tilde{T})^{3/2}} - \frac{\|u\|^2}{2\tilde{T}}$$

$$\alpha_2 = \frac{u}{\tilde{T}} \quad \alpha_3 = -\frac{1}{\tilde{T}}$$

THE NON LINEAR SYSTEM FOR $\vec{\alpha}$ IS:

$$\left\| \sum_{k=1}^N \vec{m}(v_k) \exp(\vec{\alpha} \cdot \vec{m}(v_k)) \omega_k = \begin{pmatrix} \rho \\ pu \\ E \end{pmatrix} \right.$$

$$\Leftrightarrow \boxed{\Phi(\vec{\alpha}) = 0}$$

EQUATION : $\Phi(\vec{\alpha}) = 0$

WHERE $\Phi(\vec{\alpha}) = \sum_{k=1}^N \vec{m}(v_k) \exp(\vec{\alpha} \cdot \vec{m}(v_k)) \omega_k$
- $\begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}$

NEWTON METHOD: - INITIAL GUESS $\vec{\alpha}_0$

- LINEARIZE Φ AROUND $\vec{\alpha}_0$:

$$\Phi(\vec{\alpha}) \approx \Phi(\vec{\alpha}_0) + D\Phi(\vec{\alpha}_0) \cdot (\vec{\alpha} - \vec{\alpha}_0)$$

- $\vec{\alpha}_1$ IS SEEKED AS A ZERO OF THIS
AFFINE MODEL : $\vec{\alpha}_1$ S.T

$$D\Phi(\vec{\alpha}_0) \cdot (\vec{\alpha}_1 - \vec{\alpha}_0) = -\Phi(\vec{\alpha}_0) \quad \text{LINEAR SYSTEM } 5 \times 5$$

- ITERATE THIS : WE HAVE $\vec{\alpha}_r$, WE SOLVE

$$D\Phi(\vec{\alpha}_r) \cdot (\vec{\alpha}_{r+1} - \vec{\alpha}_r) = -\Phi(\vec{\alpha}_r)$$

UNTIL CONVERGENCE ($\|\Phi(\vec{\alpha}_r)\|$ IS SUFFICIENTLY SMALL)

INITIAL GUESS: $\vec{\alpha}_0$ TWO POSSIBILITIES

$$\textcircled{1} \quad \vec{\alpha}_0 = \left(\log \frac{\rho}{(2\pi T)^{3/2}} - \frac{\|u\|^2}{2T}, \frac{u}{T}, -\frac{1}{T} \right)$$

(THIS THE VALUE OF $\vec{\alpha}$ FOR A CONTINUOUS VELOCITY SPACE)

L) WHEN WE HAVE A "NOT TOO SMALL" NUMBER OF VELOCITIES: GOOD APPROXIMATION

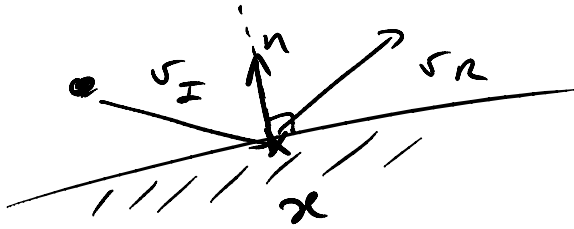
$$\textcircled{2} \quad \vec{\alpha}_0 = \vec{\alpha} \text{ COMPUTED AT THE PREVIOUS TIME ITERATION}$$

NEWTON ALG. CONVERGES VERY FAST

$$(\uparrow \leq 5)$$

L1. L1 APPROXIMATION OF BOUNDARY

COLLISIONS



$$v_R \cdot n \geq 0$$

$$f(t, x, v_R) = \mathcal{R}(f(t, x, \cdot)) \quad (*)$$

RELATION IS S.T : ZERO MASS FLUX ACROSS THE SOLID BOUNDARY

$$\int_{\mathbb{R}^3} v \cdot n f(t, x, v) dv = 0 \quad (**)$$

PROBLEM : HOW TO DISCRETIZE THE BOUNDARY CONDITION (*) SO AS TO PRESERVE (**)?

4.4.1 DIFFUSE REFLECTION

$$f(t, x, v) = M[\sigma_w, u_w, T_w](v) \quad \text{FOR } v \cdot n \geq 0$$

WITH • u_w AND T_w ARE VELOCITY AND TEMPERATURE OF THE WALL

- σ_w IS DEFINED SO AS TO ENSURE (**):

$$\sigma_w = \frac{\left(\int_{v \cdot n \leq 0} (v \cdot n) f(t, x, v) dv \right)}{\left(\int_{v \cdot n \geq 0} (v \cdot n) M[\sigma_w, u_w, T_w](v) dv \right)}$$

AT THE DISCRETE LEVEL:

- LET M^W THE DISCRETE MAXWELLIAN S.T

$$\sum_{k=1}^N \begin{pmatrix} 1 \\ v_k \\ \frac{1}{2} \|v_k\|^2 \end{pmatrix} M_k^W \omega_k = \begin{pmatrix} 1 \\ u_w \\ E_w \end{pmatrix}$$

$$(E_w = \frac{1}{2} \|u_w\|^2 + \frac{3}{2} T_w)$$

(CONSERVATIVE APPROX. OF $M[1, u_w, T_w]$ ON THE GRID V)

THE DISCRETE BOUNDARY RELATION IS:

$$F_k(t, x) = \sigma_w M_k^W \quad \text{FOR } k \text{ ST } v_{k \cdot n} \geq 0$$

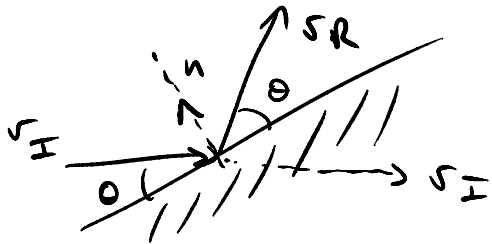
WHERE

$$\sigma_w = - \frac{\sum_{k \cdot (v_{k \cdot n} \leq 0)} (v_{k \cdot n}) F_k \omega_k}{\sum_{k \cdot (v_{k \cdot n} \geq 0)} (v_{k \cdot n}) M_k^W \omega_k}$$

EXERCISE: PROVE THAT

$$\sum_{k=1}^N (v_{k \cdot n}) F_k(t, x) \omega_k = 0 \quad (\text{ZERO MASS FLUX ACROSS THE WALL IS PRESERVED})$$

4.4.2 SPECULAR REFLECTION



SAME ANGLE FOR v_I AND v_R

$$\Leftrightarrow v_R = \text{ROTATION}_{2\theta}(v_I)$$

$$\Leftrightarrow v_I = \text{ROTATION}_{-2\theta}(v_R)$$

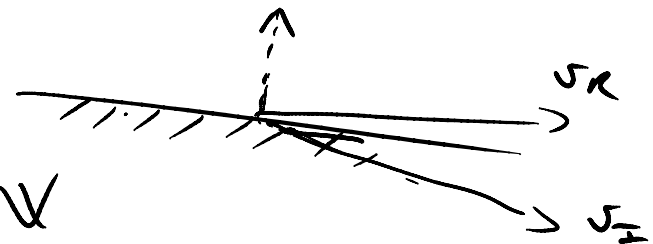
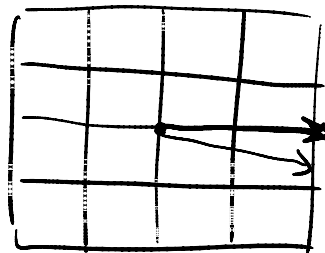
$$\left| f(t, x, v_R) = f(t, x, \text{ROTATION}_{-2\theta}(v_R)) \right|$$

DISCRETE LEVEL: THE PROBLEM IS THAT

IF $v_R \in \mathcal{V}$ THERE IS NO REASON

FOR $v_I = \text{ROTATION}_{-2\theta}(v_R) \in \mathcal{V}$

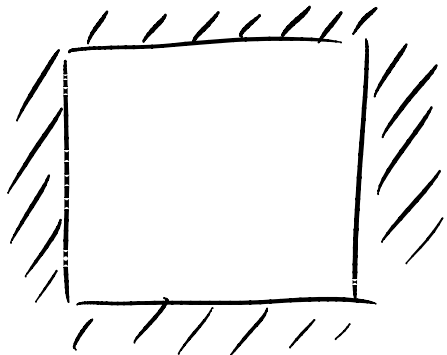
EXAMPLE (2):



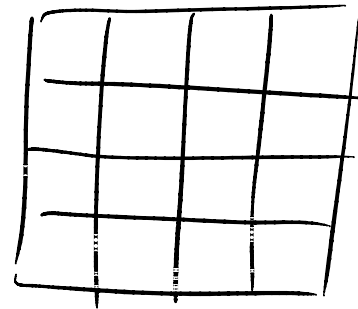
WE NEED SOME INTERPOLATION TO APPROXIMATE F AT $\mathcal{V}_I \notin \mathcal{V}$

SO AS TO PRESERVE THE ZERO MASS FLUX CONDITION.

EXCEPT IF THE GRID IS INVARIANT THROUGH THE ROTATION : EXAMPLE



PHYSICAL SPACE
(FOR x)



VELOCITY GRID \mathcal{V}

EXERCISE:

PROVE THAT : FOR EVERY $\mathcal{V}_K \in \mathcal{V}$ \mathcal{V}_I IS NECESSARILY IN \mathcal{V} TOO.