

Network-Based Kinetic Models: Emergence of a Statistical Description of the Graph Topology

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- Interacting multi-agent systems are ubiquitous in both classical physics and socio-/econophysics
- In socio-/econophysics, unlike classical physics, interactions are often networked: agents do not interact "all-to-all" but according to some preferential connections
- Prototype: opinion formation on social networks
- Large number of networked agents \rightsquigarrow need for a statistical description of the network topology

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- B. Düring, J. Franceschi, M.-T. Wolfram and M. Zanella. "Breaking consensus in kinetic opinion formation models on graphons". Preprint. 2024
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Particle Description

- Agents are understood as the vertices of a graph $\mathcal{G}=(\mathcal{I},\,\mathcal{E}),\,\mathcal{I}=\{1,\,\ldots,\,N\}$
- A representative agent $X\in\mathcal{I}$ features an opinion $V_t\in\mathcal{O}\subset\mathbb{R}$ at time $t\geq 0$
- Opinion exchange algorithm in randomly selected pairs of agents:

 $V_{t+\Delta t} = (1 - \Theta)V_t + \Theta\Psi(V_t, V_t^*)$ $V_{t+\Delta t}^* = (1 - \Theta)V_t^* + \Theta\Psi_*(V_t^*, V_t), \qquad \Theta \sim \text{Bernoulli}(B(X, X_*)\Delta t)$

in an interaction time step $0<\Delta t\leq 1$

• The interaction rate ${\cal B}$ encodes the information on agents' connections:

$$B(X, X_*) = \begin{cases} 1 & \text{if } (X, X_*) \in \mathcal{E} \\ 0 & \text{otherwise} \end{cases}$$

• $\Psi,\,\Psi_*:\mathcal{O}^2\to\mathcal{O}$ represent the post-interaction opinions in case of a successful interaction

• Let
$$(X, V_t) \sim f(x, v, t)$$
, $x \in \mathcal{I}$, $v \in \mathcal{O}$

$$f(x, v, t) = \frac{1}{N} \sum_{i \in \mathcal{I}} f_i(v, t) \otimes \delta(x - i)$$

with $f_i:\mathcal{O}\times[0,+\infty)\to\mathbb{R}_+$ the pdf of the opinion of agent i at time t

• Taking the expectation of $\Phi(X, V_{t+\Delta t})$ and of $\Phi(X_*, V_{t+\Delta t}^*)$, where Φ is an arbitrary scalar function, one obtains

$$\frac{d}{dt} \sum_{h \in \mathcal{I}} \int_{\mathcal{O}} \Phi(h, v) f_h(v, t) \, dv = \\ = \sum_{h, k \in \mathcal{I}} \iint_{\mathcal{O}^2} B(h, k) \frac{\Phi(h, v') + \Phi(k, v'_*) - \Phi(h, v) - \Phi(k, v_*)}{2N} f_h(v, t) f_k(v_*, t) \, dv \, dv_*$$

where

$$v' = \Psi(v, v_*), \qquad v'_* = \Psi_*(v_*, v)$$

Derivation of a Kinetic Description 2/2

• Taking $\Phi(x,v)=\phi(x)\varphi(v)$ with

$$\phi(x) = \begin{cases} 1 & \text{if } x = i \in \mathcal{I} \\ 0 & \text{otherwise,} \end{cases} \qquad \varphi : \mathcal{O} \to \mathbb{R} \text{ arbitrary}$$

we get

$$\begin{split} \frac{d}{dt} \int_{\mathcal{O}} \varphi(v) f_i(v,t) \, dv &= \frac{1}{2N} \sum_{k \in \mathcal{I}} B(i,k) \iint_{\mathcal{O}^2} (\varphi(v') - \varphi(v)) f_i(v,t) f_k(v_*,t) \, dv \, dv_* \\ &+ \frac{1}{2N} \sum_{h \in \mathcal{I}} B(h,i) \iint_{\mathcal{O}^2} (\varphi(v'_*) - \varphi(v)) f_h(v,t) f_i(v_*,t) \, dv \, dv_* \end{split}$$

• Introducing the adjacency matrix $\mathbf{M} := (B(i, j))_{i, j \in \mathcal{I}} \in \mathbb{R}^{N \times N}$ of \mathcal{G} :

$$\frac{d}{dt} \int_{\mathcal{O}} \varphi(v) \mathbf{f}(v,t) \, dv = \frac{1}{2N} \iint_{\mathcal{O}^2} (\varphi(v') - \varphi(v)) \mathbf{f}(v,t) \odot \mathbf{M} \mathbf{f}(v_*,t) \, dv \, dv_* \\ + \frac{1}{2N} \iint_{\mathcal{O}^2} (\varphi(v'_*) - \varphi(v)) \mathbf{M}^T \mathbf{f}(v,t) \odot \mathbf{f}(v_*,t) \, dv \, dv_*$$

with $\mathbf{f}(v,t) = (f_i(v,t))_{i \in \mathcal{I}}$ and $\odot =$ Hadamard's vector product

• The system of kinetic equations

$$\frac{d}{dt} \int_{\mathcal{O}} \varphi(v) \mathbf{f}(v,t) \, dv = \frac{1}{2N} \iint_{\mathcal{O}^2} (\varphi(v') - \varphi(v)) \mathbf{f}(v,t) \odot \mathbf{M} \mathbf{f}(v_*,t) \, dv \, dv_* \\ + \frac{1}{2N} \iint_{\mathcal{O}^2} (\varphi(v'_*) - \varphi(v)) \mathbf{M}^T \mathbf{f}(v,t) \odot \mathbf{f}(v_*,t) \, dv \, dv_*$$

for the array ${\bf f}$ of opinion distribution functions is valid on whatever ${\bf graph}$

• **Problem**: it requires a "pointwise" description of the graph connections, which gets readily unfeasible when the size N of G grows

• Global opinion distribution (*v*-marginal of *f*):

$$F(v,t) := \int_{\mathcal{I}} f(x,v,t) \, dx = \frac{1}{N} \sum_{i \in \mathcal{I}} f_i(v,t) = \frac{1}{N} \mathbf{1}^T \mathbf{f}(v,t)$$

with $\mathbf{1}=(1,\,\ldots,\,1)\in\mathbb{R}^N$

• The equation for F is not closed in general:

$$\frac{d}{dt} \int_{\mathcal{O}} \varphi(v) F(v,t) \, dv =$$

= $\frac{1}{2N^2} \iint_{\mathcal{O}^2} (\varphi(v') + \varphi(v'_*) - \varphi(v) - \varphi(v_*)) \mathbf{f}^T(v,t) \mathbf{M} \mathbf{f}(v_*,t) \, dv \, dv.$

• A preliminary idea is to see whether this equation closes at least for special classes of interaction rules

Polarised Memory Interactions

- We say that an interaction rule v' = Ψ(v, v_{*}) is of polarised memory type if Ψ depends only on either v or v_{*}
 - If $v' = \Psi(v)$ we say that the interaction rule has perfect memory
 - If $v' = \Psi(v_*)$ we say that the interaction rule is memoryless
- To fix the ideas, in the following we will focus on the case

$$v' = \Psi(v), \qquad v'_* = \Psi_*(v),$$

- i.e. v' has perfect memory whereas v'_* is memoryless
- In this case:

$$\begin{aligned} &\frac{d}{dt} \int_{\mathcal{O}} \varphi(v) F(v,t) \, dv = \\ &= \frac{1}{N^2} \int_{\mathcal{O}} \left((\mathbf{w}^+)^T \frac{\varphi(v') + \varphi(v'_*)}{2} - \frac{(\mathbf{w}^-)^T + (\mathbf{w}^+)^T}{2} \varphi(v) \right) \mathbf{f}(v,t) \, dv \end{aligned}$$

with w^- , w^+ vectors of **incoming** and **outgoing degrees** of the vertices of \mathcal{G} • Notice: information about M is *lumped* in w^- , w^+

Statistical Distribution of the Degrees

 To obtain a kinetic formulation free from references to single vertices we augment the space of microscopic states by including also information on the connections via the incoming and outgoing degrees:

$$g_N(v, w^-, w^+, t) := \frac{1}{N} \sum_{\substack{i \in \mathcal{I} \\ \text{indeg}(i) = w^- \\ \text{outdeg}(i) = w^+}} f_i(v, t), \qquad w^-, w^+ \in \{0, \dots, N\}$$

• Then:

$$F(v,t) = \sum_{w^-,w^+=0}^{N} g_N(v,w^-,w^+,t), \qquad (\mathbf{w}^{\pm})^T \mathbf{f}(v,t) = N \sum_{w^-,w^+=0}^{N} w^{\pm} g_N(v,w^-,w^+,t)$$

whence we deduce a **closed equation** for g_N :

$$\frac{d}{dt} \sum_{w^-, w^+=0}^N \int_{\mathcal{O}} \varphi(v) g_N(v, w^-, w^+, t) \, dv =$$

= $\frac{1}{N} \sum_{w^-, w^+=0}^N \int_{\mathcal{O}} \left(w^+ \frac{\varphi(v') + \varphi(v'_*)}{2} - \frac{w^- + w^+}{2} \varphi(v) \right) g_N(v, w^-, w^+, t) \, dv$

Formal Limit of Growing Graph $(N o \infty)$ 1/2

• Scaling:

$$\tilde{w}^{\pm} := \frac{w^{\pm}}{N} \in \mathcal{W}_N := \left\{ \frac{n}{N}, \, n = 0, \, \dots, \, N \right\}, \qquad \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) := N^2 g_N(v, N\tilde{w}^-, N\tilde{w}^+, t)$$

• Introduce the steps $\Delta \tilde{w}^{\pm} := \frac{1}{N}$ so that

$$\sum_{\tilde{w}^{-}, \tilde{w}^{+} \in \mathcal{W}_{N}} \int_{\mathcal{O}} \tilde{g}(v, \tilde{w}^{-}, \tilde{w}^{+}, t) \, dv \, \Delta \tilde{w}^{-} \, \Delta \tilde{w}^{+} = \sum_{w^{-}, w^{+} = 0}^{N} \int_{\mathcal{O}} g_{N}(v, w^{-}, w^{+}, t) \, dv = 1$$

and the r.h.s. may be understood as a **Riemann sum** approximating, for every N, the integral of the pdf $\int_{\mathcal{O}} \tilde{g} dv$ on the square mesh of $[0, 1]^2$ produced by the grid $\mathcal{W}_N \times \mathcal{W}_N \rightsquigarrow$ cf. a graphon

• Moreover:

$$\frac{d}{dt} \sum_{\tilde{w}^-, \tilde{w}^+ \in \mathcal{W}_N} \int_{\mathcal{O}} \varphi(v) \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \, dv \, \Delta \tilde{w}^- \, \Delta \tilde{w}^+ = \\
= \sum_{\tilde{w}^-, \tilde{w}^+ \in \mathcal{W}_N} \int_{\mathcal{O}} \left(\tilde{w}^+ \frac{\varphi(v') + \varphi(v'_*)}{2} - \frac{\tilde{w}^- + \tilde{w}^+}{2} \varphi(v) \right) \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \, dv \, \Delta \tilde{w}^- \, \Delta \tilde{w}^+$$

• Passing formally to the limit $N \to \infty$, the Riemann sums w.r.t. \tilde{w}^{\pm} become integrals:

$$\begin{split} &\frac{d}{dt} \iint_{[0,\,1]^2} \int_{\mathcal{O}} \varphi(v) \tilde{g}(v,\tilde{w}^-,\tilde{w}^+,t) \, dv \, d\tilde{w}^- \, d\tilde{w}^+ = \\ &= \iint_{[0,\,1]^2} \int_{\mathcal{O}} \left(\tilde{w}^+ \frac{\varphi(v') + \varphi(v'_*)}{2} - \frac{\tilde{w}^- + \tilde{w}^+}{2} \varphi(v) \right) \tilde{g}(v,\tilde{w}^-,\tilde{w}^+,t) \, dv \, d\tilde{w}^- \, d\tilde{w}^+ \end{split}$$

• This is a single kinetic equation in which the **pointwise information** on the graph topology encoded in M has been **replaced** asymptotically by the **statistical distribution** of the (normalised) incoming and outgoing degrees of the vertices

• For general interaction rules:

$$v' = \Psi(v, v_*), \qquad v'_* = \Psi_*(v, v_*)$$

the kinetic equation for F can be written, using \tilde{g} , as:

$$\frac{d}{dt} \sum_{\tilde{w}^-, \tilde{w}^+ \in \mathcal{W}_N} \int_{\mathcal{O}} \varphi(v) \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \, dv \, \Delta \tilde{w}^- \, \Delta \tilde{w}^+ = \\
= \frac{1}{2N^2} \iint_{\mathcal{O}^2} \left(\varphi(v') + \varphi(v'_*) \right) \mathbf{f}^T(v, t) \mathbf{M} \mathbf{f}(v_*, t) \, dv \, dv_* \\
- \frac{1}{2} \sum_{\tilde{w}^-, \tilde{w}^+ \in \mathcal{W}_N} \int_{\mathcal{O}} \left(\tilde{w}^- + \tilde{w}^+ \right) \varphi(v) \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \, dv \, \Delta \tilde{w}^- \, \Delta \tilde{w}^+$$

• The first term on the r.h.s. cannot be closed in terms of the statistics of the graph connections only. In general, it requires a pointwise knowledge of the connections

$$\mathbf{M} \approx \frac{\mathbf{w}^+(\mathbf{w}^-)^T}{M_N}, \qquad M_N := \sum_{i \in \mathcal{I}} \operatorname{indeg}(i) = \sum_{i \in \mathcal{I}} \operatorname{outdeg}(i)$$

is a natural rank-one approximation of ${\bf M}$ with given incoming/outgoing degrees

• Within this approximation it results:

$$\frac{1}{2N^2} \iint_{\mathcal{O}^2} \left(\varphi(v') + \varphi(v'_*) \right) \mathbf{f}^T(v, t) \mathbf{M} \mathbf{f}(v_*, t) \, dv \, dv_* \approx \\ \approx \frac{N^2}{2M_N} \sum_{\tilde{w}^*_*, \tilde{w}^+_* \in \mathcal{W}_N} \sum_{\tilde{w}^-, \tilde{w}^+ \in \mathcal{W}_N} \iint_{\mathcal{O}^2} \tilde{w}^+ \tilde{w}^-_* \left(\varphi(v') + \varphi(v'_*) \right) \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \\ \times \tilde{g}(v_*, \tilde{w}^-_*, \tilde{w}^+_*, t) \, dv \, dv_* \, \Delta \tilde{w}^- \, \dots \, \Delta \tilde{w}^+_*$$

• Moreover, it can be shown that

$$\frac{M_N}{N^2} \xrightarrow{N \to \infty} m := \iint_{[0, 1]^2} \int_{\mathcal{O}} \tilde{w}^{\pm} \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \, dv \, d\tilde{w}^- \, d\tilde{w}^+$$

General Formal Limit of Growing Graph $(N \to \infty)$

• Within the rank-one approximation of \mathbf{M} , the equation for \tilde{g} converges formally to:

$$\frac{d}{dt} \iint_{[0, 1]^2} \int_{\mathcal{O}} \varphi(v) \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \, dv \, d\tilde{w}^- \, d\tilde{w}^+ = \\
= \frac{1}{2} \iiint_{[0, 1]^4} \iint_{\mathcal{O}^2} \frac{\tilde{w}^+ \tilde{w}^-_*}{m} \big(\varphi(v') + \varphi(v'_*) - \varphi(v) - \varphi(v_*) \big) \\
\times \tilde{g}(v, \tilde{w}^-, \tilde{w}^+, t) \tilde{g}(v_*, \tilde{w}^-_*, \tilde{w}^+_*, t) \, dv \, dv_* \, d\tilde{w}^- \, \dots \, d\tilde{w}^+_*$$

• Interestingly, this is a **classical Boltzmann-type equation** for the distribution function \tilde{g} on the space state $\mathcal{O} \times [0, 1]^2$ with

$$\frac{\tilde{w}^+\tilde{w}^-_*}{m}$$

as collision kernel