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SPECTRAL DEVIATIONS FOR THE DAMPED WAVE EQUATION

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Abstract. We prove a Weyl-type fractal upper bound for the spectrum of the damped wave equation, on a negatively curved compact manifold. It is known that most of the eigenvalues have an imaginary part close to the average of the damping function. We count the number of eigenvalues in a given horizontal strip deviating from this typical behaviour; the exponent that appears naturally is the “entropy” that gives the deviation rate from the Birkhoff ergodic theorem for the geodesic flow. A Weyl-type lower bound is still far from reach; but in the particular case of arithmetic surfaces, and for a strong enough damping, we can use the trace formula to prove a result going in this direction.

1 Results and Questions about the Spectrum of the Damped Wave Equation

Let (M, g) be a smooth compact Riemannian manifold without boundary. Let a be a C^∞ real valued function on M . We study the “damped wave equation”,

$$(\partial_t^2 - \Delta + 2a(x)\partial_t)v = 0 \tag{1.1}$$

for $t \in \mathbb{R}$ and $x \in M$. (The term “damped” applies to the case when $a \geq 0$, that is to say, when the energy is decreasing. In this paper the sign of a will be of no importance.) We shall be interested in the stationary solutions, that is, solutions of the form $v(t, x) = e^{it\tau}u(x)$ for some $\tau \in \mathbb{C}$. This means that u must satisfy

$$(-\Delta - \tau^2 + 2ia\tau)u = 0. \tag{1.2}$$

Equivalently, τ is an eigenvalue of the operator

$$\begin{pmatrix} 0 & I \\ -\Delta & 2ia \end{pmatrix}$$

acting on $H^1(M) \times L^2(M)$. For $a = 0$ this operator is the (anti-selfadjoint) wave operator; but for $a \neq 0$ the operator is not normal anymore. It is known that its spectrum is discrete and consists of a sequence (τ_n) with $\Im m(\tau_n)$ bounded and $|\Re e(\tau_n)| \rightarrow +\infty$ (see section 2). One sees easily that $\Im m(\tau_n) \in [2 \min(\inf a, 0), 2 \max(\sup a, 0)]$ if $\Re e(\tau_n) = 0$, and $\Im m(\tau_n) \in [\inf a, \sup a]$ if $\Re e(\tau_n) \neq 0$ [Sj]. One can also note, obviously, that $-\bar{\tau}_n$ is an eigenvalue if τ_n is; the spectrum is symmetric with respect to the imaginary axis.

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1.1 Background. Motivated by questions from control theory, Lebeau [Le] was interested in the so-called “stabilization problem”. Define

$$\rho = \sup \{ \beta, \exists C, E(u_t) \leq C e^{-\beta t} E(u_0) \text{ for every solution } u \text{ of (1.1)} \}$$

where the energy functional is $E(u) = \int_M |\nabla u|^2$. The stabilization problem consists in finding some damping function a (necessarily non-negative) such that $\rho > 0$. Lebeau identified

$$\rho = 2 \min \{ D(0), C(\infty) \}$$

where $D(0)$ is a sort of “spectral gap”:

$$D(0) = \inf \{ \Im m(\tau_n), \tau_n \neq 0 \}$$

and $C(\infty)$ is defined in terms of the Birkhoff averages of a along the geodesic flow

$$G^t : T^*M \longrightarrow T^*M,$$

$$C(\infty) = \lim_{t \rightarrow +\infty} \inf_{\rho \in S^*M} \frac{1}{t} \int_0^t a(G^s \rho) ds.$$

Lebeau also showed, on an example, that it is possible to have a spectral gap ($D(0) > 0$) but no exponential damping ($\rho = 0$); this surprising phenomenon is typical of non-normal operators.

Markus and Matsaev [MM] proved an analogue of Weyl’s law, also found independently later on by Sjöstrand [Sj]:

$$\#\{n, 0 \leq \Re e(\tau_n) \leq \lambda\} = \left(\frac{\lambda}{2\pi}\right)^d \left(\int_{p^{-1}[0, \lambda]} dx d\xi + \mathcal{O}(\lambda^{-1}) \right) \tag{1.3}$$

where d is the dimension of M , p is the principal symbol of $-\Delta$, namely the function $p(x, \xi) = g_x(\xi, \xi)$ defined on the cotangent bundle T^*M , and $dx d\xi$ is the Liouville measure on T^*M (coming from its canonical symplectic structure).

It is a natural question to ask about the distribution of the imaginary parts $\Im m(\tau_n)$ in the interval $[\inf a, \sup a]$. For non-normal operators, obtaining fine information on the distribution of the spectrum is much harder than for normal operators – one reason being the absence of continuous (or even smooth) functional calculus. Another related difficulty is that there is no general relation between the norm of the resolvent and the distance to the spectrum. Some techniques have been developed to obtain upper bounds on the density of eigenvalues, but lower bounds are notoriously more difficult.

Assuming that the geodesic flow is ergodic with respect to the Liouville measure on an energy layer, Sjöstrand proved that most of the τ_n have asymptotically an imaginary part $\Im m(\tau_n) \sim \bar{a}$, where \bar{a} denotes the average of a on the energy layer $S^*M = \{(x, \xi) \in T^*M, g_x(\xi, \xi) = 1\}$ with respect to the Liouville measure.

Theorem 1.1 [Sj]. *Assume that the geodesic flow is ergodic with respect to the Liouville measure on S^*M . For every $C > 0$, for every $\epsilon > 0$, we have*

$$\#\{n, \lambda \leq \Re e(\tau_n) \leq \lambda + C, \Im m(\tau_n) \notin [\bar{a} - \epsilon, \bar{a} + \epsilon]\} = o(\lambda^{d-1})$$

as λ goes to infinity.

REMARK 1.2. If C is not too small, one sees from (1.3) that there exists $\tilde{C} > 0$ such that

$$\#\{n, \lambda \leq \Re e(\tau_n) \leq \lambda + C\} \geq \tilde{C}\lambda^{d-1}$$

for large λ . Thus, the theorem does say that a majority of the τ_n have $\Im m(\tau_n) \in [\bar{a} - \epsilon, \bar{a} + \epsilon]$.

REMARK 1.3. More generally, without the ergodicity assumption, [Sj] proves the following results. Introduce the functions on T^*M ,

$$\begin{aligned} \langle a \rangle_T &= \frac{1}{T} \int_{-T/2}^{T/2} a \circ G^s ds, \\ \langle a \rangle_\infty &= \lim_{T \rightarrow +\infty} \langle a \rangle_T, \end{aligned}$$

and the real numbers

$$\begin{aligned} a_+ &= \lim_{T \rightarrow +\infty} \sup_{S^*M} \langle a \rangle_T, \\ a_- &= \lim_{T \rightarrow +\infty} \inf_{S^*M} \langle a \rangle_T. \end{aligned}$$

The function $\langle a \rangle_\infty$ is well defined Liouville-almost-everywhere, by the Birkhoff theorem.

Lebeau [Le] proves that, for any $\epsilon > 0$, there are at most finitely many n with $\Im m(\tau_n) \notin [a_- - \epsilon, a_+ + \epsilon]$, and Sjöstrand [Sj] proves that

$$\#\{n, \lambda \leq \Re e(\tau_n) \leq \lambda + C, \Im m(\tau_n) \notin [ess \inf \langle a \rangle_\infty - \epsilon, ess \sup \langle a \rangle_\infty + \epsilon]\} = o(\lambda^{d-1}).$$

1.2 Semiclassical formulation. As in [Sj] we reformulate the problem using semiclassical notation. In doing so, we also generalize the problem a little. We introduce a semiclassical parameter $0 < \hbar \ll 1$ and will be interested in the eigenvalues τ such that $\hbar\tau \xrightarrow{\hbar \rightarrow 0} 1$. If we put $\tau = \lambda/\hbar$ with λ close to 1, then equation (1.2) can be rewritten as

$$\left(-\frac{\hbar^2 \Delta}{2} - \frac{\lambda^2}{2} + i\hbar\lambda a \right) u = 0, \tag{1.4}$$

or

$$(\mathcal{P} - z)u = 0 \tag{1.5}$$

if we put $z = \lambda^2/2$, and

$$\mathcal{P} = \mathcal{P}(z) = P + i\hbar Q(z), \quad P = -\frac{\hbar^2 \Delta}{2}, \quad Q(z) = a\sqrt{z}. \tag{1.6}$$

The parameter z is close to $E = 1/2$.

More generally, we will consider a spectral problem of the form (1.5) where

$$\mathcal{P} = \mathcal{P}(z) = P + i\hbar Q(z), \quad P = -\frac{\hbar^2 \Delta}{2},$$

$z \in \Omega = e^{i] - s_0, s_0[} E_{\min}, E_{\max}[$, with $0 < E_{\min} < 1/2 < E_{\max} < +\infty$, $0 < s_0 < \pi/4$. We will assume that $Q(z) \in \Psi DO^1$ is a pseudodifferential operator that depends holomorphically on $z \in \Omega$, and that Q is formally self-adjoint for z real (the definition of our operator classes is given in section 10).

We denote

$$\Sigma = \Sigma_{\hbar} = \{z \in \Omega, \exists u, (\mathcal{P}(z) - z)u = 0\}.$$

The operator P has principal symbol $p_o(x, \xi) = g_x(\xi, \xi)/2$, and $Q(z)$ has principal symbol $q_z(x, \xi)$, taking real values for z real. In this notation, the previous results read as follows for any $E_{\min} < E_1 \leq E_2 < E_{\max}$:

$$\#\{z \in \Sigma, E_1 \leq \Re e(z) \leq E_2\} = \frac{1}{(2\pi\hbar)^d} \left[\int_{p_o^{-1}[E_1, E_2]} dx d\xi + \mathcal{O}(\hbar) \right] \quad (1.7)$$

(uniformly for all $c > 0$ and for all E_1, E_2 such that $|E_2 - E_1| \geq 2c\hbar$, E_1 and E_2 staying away from E_{\min}, E_{\max}). One can show that $\Im m(z)/\hbar$ (for $z \in \Sigma$) has to stay bounded if E_1, E_2 stay in a bounded interval. Fix some $c > 0$, and take $E_1 = E - c\hbar$ and $E_2 = E + c\hbar$. Let us denote

$$q_E^- = \lim_{T \rightarrow +\infty} \inf_{p_o^{-1}\{E\}} \langle q_E \rangle_T,$$

$$q_E^+ = \lim_{T \rightarrow +\infty} \sup_{p_o^{-1}\{E\}} \langle q_E \rangle_T,$$

then we have, in the limit $\hbar \rightarrow 0$ ([Sj]),

$$q_E^- + o(1) \leq \frac{\Im m(z)}{\hbar} \leq q_E^+ + o(1), \quad (1.8)$$

for all $z \in \Sigma$, such that $E - c\hbar \leq \Re e(z) \leq E + c\hbar$. Finally, denote \bar{q}_E the average of q_E on the energy layer $p_o^{-1}\{E\}$. Assuming that the geodesic flow is ergodic on $p_o^{-1}\{E\}$, then for any $\epsilon > 0$, any $c > 0$,

$$\#\left\{z \in \Sigma, E - c\hbar \leq \Re e(z) \leq E + c\hbar, \frac{\Im m(z)}{\hbar} \notin [\bar{q}_E - \epsilon, \bar{q}_E + \epsilon]\right\} = o(\hbar^{1-d}). \quad (1.9)$$

The method of [Sj] allows us to treat the case of a more general P (and thus a more general Hamiltonian flow than the geodesic flow), and also to deal with the case when the flow is not ergodic. However, in this paper we will stick to the case of an ergodic geodesic flow; we will add even stronger assumptions in the next paragraph.

1.3 Questions, and a few results. We try to give (partial) answers to the three natural questions:

(Q1) (Fractal Weyl law) Let I be an interval that does not contain \bar{q}_E . Can we describe in a finer way the asymptotic behaviour for

$$\#\left\{z, E - c\hbar \leq \Re e(z) \leq E + c\hbar, \frac{\Im m(z)}{\hbar} \in I\right\}?$$

For instance, can we find

$$\lim_{\hbar \rightarrow 0} \frac{\log \#\left\{z \in \Sigma, E - c\hbar \leq \Re e(z) \leq E + c\hbar, \frac{\Im m(z)}{\hbar} \in I\right\}}{\log \hbar}?$$

(Q2) (Inverse problem) Suppose we know everything about P and about the dynamics of the geodesic flow, but that Q is unknown. To what extent does the knowledge of the eigenfrequencies $\{z \in \Sigma, E - c\hbar \leq \Re e(z) \leq E + c\hbar\}$ determine q_E ?

Replacing \mathcal{P} by $B^{-1}\mathcal{P}B$, where B is an elliptic pseudodifferential operator with positive principal symbol b , amounts to replacing q by $q - \{p_o, \log b\}$, where $\{\cdot, \cdot\}$ stands for the Poisson bracket on T^*M . Two smooth functions f and g on T^*M are said to be cohomologous (with respect to the geodesic flow) if there exists a third smooth function h such that $f = g + \{p_o, h\}$. This defines an equivalence relation, we write $f \sim_{p_o} g$.

It is thus more natural to ask:

(Q2') Does the knowledge of the eigenfrequencies $\{z \in \Sigma, E - c\hbar \leq \Re e(z) \leq E + c\hbar\}$ determine the *cohomology class* of q_E ?

If the length spectrum of M is simple, then one can most probably use a trace formula and, from the knowledge of Σ , recover all the integrals of q_E along closed geodesics. And this is enough to determine the cohomology class of q_E if M has negative sectional curvature. The Livshitz theorem [Li], [GK] says that if two functions have the same integrals along all closed geodesics, then they are cohomologous. Thus, the answer to (Q2) is probably “yes” if M has negative sectional curvature and the length spectrum is simple (this last assumption is satisfied generically, but unfortunately not for surfaces of constant negative curvature).

In fact, it also makes sense to ask whether *some* knowledge of the imaginary parts alone $\{\Im m(z), z \in \Sigma, E - c\hbar \leq \Re e(z) \leq E + c\hbar\}$ allows us to recover *some* information about q_E . For instance, one can ask the apparently simple question:

(Q3) (Very weak inverse problem) Let C denote a constant function. As follows from (1.8), we have the implication

$$q_E \sim_{p_o} C \quad \text{on} \quad p_o^{-1}\{E\} \implies \frac{\Im m(z)}{\hbar} \xrightarrow{\hbar \rightarrow 0, z \in \Sigma, E - c\hbar \leq \Re e(z) + c\hbar} C.$$

Does the converse hold?

Main assumption on the manifold M . From now on, we assume that M has constant sectional curvature -1 . This implies the ergodicity of the geodesic flow on any non-degenerate energy layer (with respect to the Liouville measure), and in fact a very chaotic behaviour of trajectories (see section 3). We will indicate how to generalize our results in the case of *surfaces* of variable negative curvature; however, it is not clear what to do in higher dimension and variable negative curvature.

In the following theorem, h_{KS} stands for the Kolmogorov–Sinai entropy, or metric entropy. It is an affine functional, taking non-negative values, defined on \mathcal{M} , the set of G -invariant probability measures on T^*M : see for instance [KH] for the definition of h_{KS} . We will also denote $\mathcal{M}_E \subset \mathcal{M}$ the set of invariant probability measures carried by $p_o^{-1}\{E\}$. Since p_o is homogeneous it is enough to consider, for instance, the case $E = 1/2$, and $p_o^{-1}\{E\} = S^*M$. For $\mu \in \mathcal{M}_{1/2}$, we have $h_{KS}(\mu) \leq d - 1$, with equality if and only if μ is the Liouville measure. We now fix $E = 1/2$ and we denote $q = q_{1/2}$, $\bar{q} = \bar{q}_{1/2}$, $q^+ = q_{1/2}^+$, $q^- = q_{1/2}^-$.

We fix some $c > 0$ and denote

$$\Sigma_{1/2} = \left\{ z \in \Sigma, \frac{1}{2} - c\hbar \leq \Re e(z) \leq \frac{1}{2} + c\hbar \right\}.$$

Theorem 1.4. *Assume M has constant sectional curvature -1 . Define*

$$H(\alpha) = \sup \left\{ h_{KS}(\mu), \mu \in \mathcal{M}_{1/2}, \int q d\mu = \alpha \right\}.$$

If $\alpha \geq \bar{q}$, then for any $c > 0$

$$\limsup_{\hbar \rightarrow 0} \frac{\log \#\{z \in \Sigma_{1/2}, \frac{\Im m(z)}{\hbar} \geq \alpha\}}{|\log \hbar|} \leq H(\alpha).$$

If $\alpha \leq \bar{q}$, then for any $c > 0$

$$\limsup_{\hbar \rightarrow 0} \frac{\log \#\{z \in \Sigma_{1/2}, \frac{\Im m(z)}{\hbar} \leq \alpha\}}{|\log \hbar|} \leq H(\alpha).$$

REMARK 1.5 (See [L2, §4], or [BL, §3], for the argument). The function H is concave and is identically $-\infty$ outside $[q^-, q^+]$. It is continuous and strictly concave in $[q^-, q^+]$, and real analytic in $]q^-, q^+[$ (note that $q^- = q^+$ if and only if q is cohomologous to a constant on S^*M). The function H reaches its maximum $d - 1$ at the point \bar{q} , and has finite (non-negative) limits at the endpoints q^-, q^+ . In particular H is positive in the open interval $]q^-, q^+[$.

REMARK 1.6. The key fact in the proof of Theorem 1.4 is the large deviation estimate proved by Kifer [Kil],

$$\lim_{T \rightarrow +\infty} \frac{\log L_{1/2}\{\rho \in S^*M, \langle q \rangle_T(\rho) \in I\}}{T} = \sup\{H(\alpha), \alpha \in I\} - d - 1$$

for any interval $I \subset \mathbb{R}$ – where $L_{1/2}$ is the Liouville measure on $p_o^{-1}\{1/2\} = S^*M$. This result gives the volume of the set of trajectories deviating from the Birkhoff ergodic theorem (see section 3). Kifer’s large deviation estimate is used in Corollary 3.

On a surface of variable negative curvature, we can generalize the result to

Theorem 1.7. *Assume M is a surface of variable negative curvature. Denote φ the infinitesimal unstable Jacobian (see section 3). Define*

$$H(\alpha) = \sup \left\{ \frac{h_{KS}(\mu)}{\int \varphi d\mu}, \mu \in \mathcal{M}_{1/2}, \int q d\mu = \alpha \right\}.$$

Then the same conclusion as in Theorem 1.4 holds.

REMARK 1.8. For a manifold of variable negative curvature and dimension d we would expect the same to hold with $H(\alpha) = (d - 1) \sup\{h_{KS}(\mu)/\int \varphi d\mu, \mu \in \mathcal{M}_{1/2}, \int q d\mu = \alpha\}$. However, our proof does not work in this case. We stress again the fact that the function H attains its maximum $d - 1$ only for $\alpha = \bar{q}$.

One may wonder if the lim sup in Theorem 1.4 is also a lim inf. This question is reminiscent of the fractal Weyl law conjecture asked by Zworski and Sjöstrand, but in our situation we can say with certainty that it is not true. Worse, one cannot expect the lower bound to hold for a “generic” q . In [S], Emmanuel Schenck proved that there exists $\delta > 0$ such that $\Im m(z)/\hbar \leq q^+ - \delta$ for all $z \in \Sigma_{1/2}$, provided a certain criterion on q is satisfied. The criterion reads $\Pr(q - \frac{d-1}{2}) < q^+$, where the pressure functional \Pr is defined on the space of continuous functions on S^*M and is

the Legendre transform of $-h_{KS}$ (see section 3). The functional Pr is lipschitz with respect to the C^0 norm, and the condition $\text{Pr}(q - \frac{d-1}{2}) < q^+$ defines a non-empty open set in the C^0 topology. For such a q we cannot have

$$\liminf_{\hbar \rightarrow 0} \frac{\log \#\{z \in \Sigma_{1/2}, \frac{\Im m(z)}{\hbar} > q_+ - \delta\}}{|\log \hbar|} \geq H(q_+ - \delta),$$

since H is positive in $]q_-, q_+[$ but the \liminf on the left-hand side is $-\infty$.

In sections 8 and 9, we investigate (Q3) in some special cases. We work on compact hyperbolic surfaces ($d = 2$), and we study operators of the form

$$\Delta_\omega f = \Delta f - 2\langle \omega, df \rangle + \|\omega\|_x^2 f,$$

called “twisted laplacians” – here ω is a harmonic real-valued 1-form on M . Studying the large eigenvalues of Δ_ω amounts to studying a fixed spectral window for the semiclassical twisted laplacian

$$-\hbar^2 \frac{\Delta_\omega}{2} = -\hbar^2 \frac{\Delta}{2} + \hbar^2 \langle \omega, d. \rangle - \hbar^2 \frac{\|\omega\|_x^2}{2}, \quad \hbar \rightarrow 0.$$

This question falls exactly into the setting of section 1.2, with $q(x, \xi) = \langle \omega_x, \xi \rangle$. Since $q(x, -\xi) = -q(x, \xi)$, we have $\bar{q} = 0$. We will denote $\text{Pr}(\omega) = \text{Pr}(q)$ the pressure of the function q , defined in section 3.2. It will also be interesting to note that $q^+ = -q^-$ coincides with the stable norm $\|\omega\|_s$ defined in section 8.2, that $\|\omega\|_s + 1 \geq \text{Pr}(\omega) \geq \|\omega\|_s$ and that $\text{Pr}(t\omega) - |t|\|\omega\|_s \xrightarrow{t \rightarrow \infty} 0$ in the case of surfaces.

Theorem 1.9. *Assume M is a compact arithmetic surface coming from a quaternion algebra. Take $\omega \neq 0$. Fix $\beta \in (0, 1]$, and $0 < \epsilon < \beta$. Then, for every \hbar small enough, there exists $z \in Sp(-\hbar^2 \frac{\Delta_\omega}{2})$ with $|\Re e(z) - \frac{1}{2}| \leq \hbar^{1-\beta}$, such that*

$$\frac{\Im m(z)}{\hbar} \geq \text{Pr}(\omega) - \frac{3}{4} - \frac{1 + \epsilon}{2\beta}.$$

Equivalently, given $\beta \in (0, 1]$, and $0 < \epsilon < \beta$, for all T large enough, there exists r_n such that $|\Re e(r_n) - T| \leq T^\beta$ and

$$|\Im m(r_n)| \geq \text{Pr}(\omega) - \frac{3}{4} - \frac{1 + \epsilon}{2\beta},$$

where r_n is the “spectral parameter” defined by $\lambda_n = -(\frac{1}{4} + r_n^2)$.

Of course, we deduce immediately the following corollary:

COROLLARY 1. $\#\{n, |\Re e(r_n)| \leq T, |\Im m(r_n)| \geq \text{Pr}(\omega) - \frac{3}{4} - \frac{1+\epsilon}{2\beta}\} \geq T^{1-\beta}$, asymptotically as $T \rightarrow +\infty$.

Since $\text{Pr}(\omega) \rightarrow 1$ as $\omega \rightarrow 0$ (but $\text{Pr}(\omega) \rightarrow +\infty$ as $\|\omega\|_s \rightarrow +\infty$) this result is only interesting if $\|\omega\|_s$ is large enough (depending on β). Note also that if $\|\omega\|_s$ is large enough we have $\text{Pr}(\omega) < \|\omega\|_s + \frac{1}{2}$, so that the work of Emmanuel Schenck mentioned above will show that there is a strip $\{\Im m(z) \geq \|\omega\|_s - \delta\}$ ($\delta > 0$) that contains no r_n .

The arithmetic case is special because the lengths of closed geodesics are well separated: we can write a trace formula, find a lower bound on the geometric part (despite its oscillatory nature) and deduce a lower bound on the imaginary parts of eigenvalues. The ideas are inspired by Hejhal’s book [H].

REMARK 1.10. Another way of defining the twisted laplacian is to write $M = \Gamma \backslash \mathbb{H}$, where \mathbb{H} is the universal cover of M and Γ is a discrete group of isometries of \mathbb{H} ; to fix an origin $o \in \mathbb{H}$, and to write

$$\Delta_\omega f(x) = e^{\int_o^x \omega} \circ \Delta (e^{-\int_o^x \omega} f(x));$$

this operator preserves Γ -periodic functions, hence descends to M . The case where ω takes purely imaginary values is self-adjoint, and analogous to the study of “Bloch waves”. The case where ω takes real values is no longer self-adjoint.

REMARK 1.11. Our motivation for working with twisted laplacians on hyperbolic manifolds was that it is a case where the trace formula is exact. There was, *a priori*, hope of proving finer results than in cases where the trace formula is not exact (in the latter case the space of test functions to which the formula can be applied is more limited). *A posteriori*, we never use any wild test function that wouldn't be allowed in the general case. So, one can think that the same result holds for our general operator (1.6) – provided one proves a semiclassical trace formula first.

If we don't assume arithmeticity, we can only treat the following operator:

$$-\hbar^2 \frac{\Delta_{\theta(\hbar)\omega}}{2} = -\hbar^2 \frac{\Delta}{2} + \hbar^2 \theta(\hbar) \langle \omega, d. \rangle - \hbar^2 \theta(\hbar)^2 \frac{\|\omega\|_x^2}{2}, \quad \hbar \rightarrow 0,$$

where $\theta(\hbar) \geq |\log \hbar|$ and $\hbar \theta(\hbar) \rightarrow 0$. In other words the non-selfadjoint perturbation is stronger than in the previous cases.

Theorem 1.12. *Assume M is a compact hyperbolic surface. Take $\omega \neq 0$. Take any function $f(\hbar) \gg \theta(\hbar)^{3/2} \log \log \hbar^{1/2}$. Then there is a sequence $\hbar_n \rightarrow 0$ such that*

$$\sup \left\{ \frac{\Im m(z)}{\hbar_n \theta(\hbar_n)}, z \in Sp \left(-\hbar_n^2 \frac{\Delta_{\theta(\hbar_n)\omega}}{2} \right), \left| \Re z - \frac{1}{2} \right| \leq \hbar_n f(\hbar_n) \right\} \xrightarrow{n \rightarrow +\infty} \|\omega\|_s,$$

the stable norm of ω .

See section 8.2 for the definition of the stable norm. Note that with our previous notation $\|\omega\|_s = q^+ = -q^-$, and $\Pr(t\omega) - |t| \|\omega\|_s \xrightarrow[t \rightarrow \infty]{} 0$ on a hyperbolic surface.

1.4 Links with quantum unique ergodicity. (Q3) has a rather close link with the question of quantum unique ergodicity on negatively curved manifolds. To explain what we mean, let us go back to the setting of the damped wave equation (our discussion could apply as well to the generalized setting of section 1.2).

An easy first remark is that

$$\Im m(\tau) = \frac{\langle u, au \rangle}{\langle u, u \rangle}$$

for a solution of (1.2). The fact that $\Im m(\tau_n)$ has no limit as $n \rightarrow +\infty$ implies that the sequence of probability measures $|u_n(x)|^2 d\text{Vol}(x)$ has no limit either (where the u_n are normalized eigenfunctions of the damped wave operator, of eigenvalue τ_n). Quantum unique ergodicity – meaning the existence of a unique limit – is expected to hold for normal operators such as the laplacian, but not for non-self-adjoint perturbations, such as the damped wave operator.

REMARK 1.13. An interesting case, for which I have no clue, is the case of eigenfunctions of $\Delta + ia$. This case is different from the ones we considered before, because the non-self-adjoint perturbation is of smaller order. It is difficult to know what to predict here concerning quantum unique ergodicity.

Here is another remark: consider the 1-parameter family of damped wave equations,

$$(-\Delta - \tau_n^2 + 2isa\tau_n)u_n = 0, \tag{1.10}$$

where s is close to 0. In other words, we see the damped wave equation as a perturbation of the spectral problem $(-\Delta - \tau_n^2)u_n = 0$. Now, $\tau_n(s)$ and $u_n(s)$ depend on s , and assuming for simplicity that the eigenvalue is simple, we can calculate $\tau_n'(0)$,

$$\tau_n'(0) = i\langle u_n(0), au_n(0) \rangle,$$

where $u_n(0)$ is the (L^2 normalized) solution to $(-\Delta - \tau_n^2)u_n = 0$. On a negatively curved manifold, the quantum unique ergodicity conjecture predicts that $\langle u_n, au_n \rangle$ tends to \bar{a} as τ_n goes to infinity – and the Shnirelman theorem says that this is the case for a subsequence of indices n of “density 1” [Sc]. In some sense, the Shnirelman theorem may be seen as an infinitesimal version (meaning $s \rightarrow 0$) of Theorem 1.1. If the quantum unique ergodicity conjecture is true, it implies that the answer to (Q3) is “no” at the infinitesimal level: (Q3) asks whether $\Im m(\tau_n) \xrightarrow{n \rightarrow +\infty} i\bar{a}$ implies that a is constant, but quantum unique ergodicity says that we should have $\tau_n'(0) \xrightarrow{n \rightarrow +\infty} i\bar{a}$ in any case. This remark shows that a positive answer to (Q3) will be easier to prove for a large than for a small – and we experienced this phenomenon in Theorems 1.9 and 1.12 (the role of a being played by ω).

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2 Note on the Definition of the Spectrum and its Multiplicity.

By the term “spectrum”, we mean the set Σ of $z \in \Omega$ such that $\mathcal{P}(z) - z$ is not bijective $H^2(M) \rightarrow H^0(M)$. If it is bijective, then the inverse must be continuous, by the closed graph theorem.

If z is restricted to a compact subset of Ω , it is easy to see that $G(z) = I + \lambda^{-1}(\mathcal{P}(z) - z)$ is invertible for $\lambda > 0$ large enough. The inverse $G(z)^{-1}$ is then a compact operator on $H^0(M)$. Besides, one sees that $\mathcal{P}(z) - z$ is not bijective $H^2(M) \rightarrow H^0(M)$ if and only if 1 is in the spectrum of $G(z)^{-1}$, if and only if there exists $u \in H^2(M)$ such that $(\mathcal{P}(z) - z)u = 0$. This shows, in particular, that the “spectrum” is discrete and corresponds to the existence of “eigenfunctions”.

To define the multiplicity of $z_0 \in \Sigma$, we proceed the same way as in [Sj]. By the density of finite-rank operators in the space of compact operators [ReS], one shows that in a neighbourhood of z_0 there exists a finite-rank operator $K(z)$, depending holomorphically on z , such that $\mathcal{P}(z_0) - z_0 + K(z_0)$ is invertible. The multiplicity of z_0 is then defined as the order of z_0 as a zero of the holomorphic function

$$z \mapsto \det[(\mathcal{P}(z) - z + K(z))^{-1}(\mathcal{P}(z) - z)] = \det[I - (\mathcal{P}(z) - z + K(z))^{-1}K(z)].$$

It is shown in [Sj] that this definition does not depend on the choice of $K(z)$ – provided it is trace class – and coincides with other usual definitions of the multiplicity.

3 A Few Facts on the Geodesic Flow on a Negatively Curved Manifold

3.1 Anosov property. If M has negative sectional curvature, then the geodesic flow on S^*M has the Anosov property [An]. This means there are $C, \lambda > 0$ such that for each $\rho \in S^*M$, the tangent space $T_\rho(S^*M)$ splits into

$$T_\rho(S^*M) = E^u(\rho) \oplus E^s(\rho) \oplus \mathbb{R}X(\rho)$$

where

- the vector field X generates the geodesic flow G^t ;
- E^s is the stable subspace: for all $v \in E^s(\rho)$, and for $t \geq 0$, $\|DG_\rho^t.v\| \leq Ce^{-\lambda t}\|v\|$;
- E^u is the unstable subspace: for all $v \in E^u(\rho)$, and for $t \leq 0$, $\|DG_\rho^t.v\| \leq Ce^{\lambda t}\|v\|$.

If M has constant negative curvature -1 , any $\lambda < 1$ will do. We take $\lambda = 1 - \epsilon$, with ϵ arbitrarily small. One also has an upper bound $\|DG_\rho^t.v\| \leq Ce^{(1+\epsilon)|t|}\|v\|$ for any $\epsilon > 0$ and any $t \in \mathbb{R}$.

3.2 Pressure, entropy, and large deviation. The pressure is defined on $C^0(S^*M)$, as the Legendre transform of the entropy:

$$\text{Pr}(f) = \sup \left\{ h_{KS}(\mu) + \int f d\mu, \mu \in \mathcal{M}_{1/2} \right\}.$$

If f is Hölder, then the supremum is attained for a unique μ , called the equilibrium measure of f . The functional Pr is analytic on any Banach space of sufficiently regular functions – for instance, a space of Hölder functions [BR], [Rue]. Besides, the restriction of Pr to any line $\{f+tg, t \in \mathbb{R}\}$ is strictly convex, unless g is cohomologous to a constant [R]. If g is sufficiently smooth, we recall that this means that $g = \{p_o, h\} + c$ for some smooth function h and a constant c . If g is Hölder, it is better to use the integral version of the notion. If $\gamma(t)$ is a periodic trajectory of the geodesic flow on S^*M (equivalently, a closed geodesic), we denote ℓ_γ its period (equivalently, the length of the closed geodesic). We denote $d\gamma$ the measure $\int g d\gamma = \int_0^{\ell_\gamma} g(\gamma(t))dt$ on S^*M , and $d\mu_\gamma$ the probability measure $\int g d\mu_\gamma = \ell_\gamma^{-1} \int g d\gamma$. One says that g is cohomologous to the constant function c if $\int g d\gamma = c\ell_\gamma$ for all periodic trajectories

of the geodesic flow (the Livschitz theorem says that both notions are equivalent for smooth functions).

Let us now fix a smooth function q on S^*M , not cohomologous to a constant. For $\alpha \in \mathbb{R}$, define

$$H(\alpha) = \sup \left\{ h_{KS}(\mu), \mu \in \mathcal{M}_{1/2}, \int q d\mu = \alpha \right\},$$

$$P(\beta) = \Pr(\beta q) = \sup \left\{ h_{KS}(\mu) + \beta \int q d\mu, \mu \in \mathcal{M}_{1/2} \right\} = \sup_{\alpha} \alpha\beta + H(\alpha).$$

The function H is concave, continuous on the interval $[q_-, q_+]$ defined earlier:

$$q^- = \lim_{T \rightarrow +\infty} \inf_{p_\sigma^{-1}\{1/2\}} \langle q \rangle_T,$$

$$q^+ = \lim_{T \rightarrow +\infty} \sup_{p_\sigma^{-1}\{1/2\}} \langle q \rangle_T.$$

In the case of a negatively curved manifold, this definition coincides with

$$q^- = \inf \left\{ \int q d\mu, \mu \in \mathcal{M}_{1/2} \right\},$$

$$q^+ = \sup \left\{ \int q d\mu, \mu \in \mathcal{M}_{1/2} \right\}.$$

The function H is real analytic and strictly concave in $]q_-, q_+[$. The function P is real analytic, strictly convex on \mathbb{R} . Clearly, $P(\beta) \geq \beta q_+$ for $\beta \geq 0$, and it is not very difficult to show that the limit $\lim_{\beta \rightarrow +\infty} P(\beta) - \beta q_+$ exists and is non-negative (this limit is equal to $H(q^+)$). Similarly, $P(\beta) \geq \beta q_-$ for $\beta \leq 0$, the limit $\lim_{\beta \rightarrow -\infty} P(\beta) - \beta q_-$ exists and is non-negative.

The pressure and the entropy appear naturally when studying large deviations for the Birkhoff averages of the function q . Denote $J_t(\rho)$ the Jacobian of DG^t going from $E^u(\rho)$ to $E^u(G^t \rho)$. Define φ (the infinitesimal unstable Jacobian) by

$$\varphi(\rho) = \frac{dJ_t}{dt} \Big|_{t=0}(\rho).$$

On a manifold of dimension d and constant negative curvature -1 , the function φ is constant, equal to $d-1$. In general one can only say that it is a Hölder function. The function φ is not necessarily positive, but it is cohomologous to a positive function, for instance $\langle \varphi \rangle_T$ for T large enough. In what follows, we will assume without loss of generality that $\varphi > 0$.

The two following large deviation results are due to Kifer [Kil].

Theorem 3.1 [Kil, Prop. 3.2]. *Let M be a compact manifold of negative sectional curvature. Let q be a smooth function on S^*M . For $T > 0$, define the function $\langle q \rangle_T = \frac{1}{T} \int_{-T/2}^{T/2} q \circ G^s ds$ on S^*M . Denote $L_{1/2}$ the Liouville measure on S^*M .*

Then

$$\lim_{T \rightarrow +\infty} \frac{\log \int_{S^*M} e^{T\langle q \rangle_T(\rho)} dL_{1/2}(\rho)}{T} = \Pr(q - \varphi).$$

As a consequence, one also has

Theorem 3.2 [Ki1, Th. 3.4]. (i) *Let M be a compact manifold of negative sectional curvature. Let q be a smooth function on S^*M . For $T > 0$, define the function $\langle q \rangle_T = \frac{1}{T} \int_{-T/2}^{T/2} q \circ G^s ds$ on S^*M . Denote $L_{1/2}$ the Liouville measure on S^*M .*

Then, for any closed interval $I \subset \mathbb{R}$, we have

$$\limsup_{T \rightarrow +\infty} \frac{\log L_{1/2} \{ \rho \in S^*M, \langle q \rangle_T(\rho) \in I \}}{T} \leq \sup \left\{ h_{KS}(\mu) - \int \varphi d\mu, \mu \in \mathcal{M}_{1/2}, \int q d\mu \in I \right\}.$$

For any open interval $I \subset \mathbb{R}$, we have

$$\liminf_{T \rightarrow +\infty} \frac{\log L_{1/2} \{ \rho \in S^*M, \langle q \rangle_T(\rho) \in I \}}{T} \geq \sup \left\{ h_{KS}(\mu) - \int \varphi d\mu, \mu \in \mathcal{M}_{1/2}, \int q d\mu \in I \right\}.$$

(ii) *Let M be a compact manifold of dimension d , with constant sectional curvature -1 . Then (i) can be rephrased as follows. Let q be a smooth function on S^*M . For $T > 0$, define the function $\langle q \rangle_T = \frac{1}{T} \int_{-T/2}^{T/2} q \circ G^s ds$ on S^*M . Denote $L_{1/2}$ the Liouville measure on S^*M . Then, for any closed interval $I \subset \mathbb{R}$, we have*

$$\limsup_{T \rightarrow +\infty} \frac{\log L_{1/2} \{ \rho \in S^*M, \langle q \rangle_T(\rho) \in I \}}{T} \leq \sup \{ H(\alpha), \alpha \in I \} - (d - 1)$$

where H is the function $H(\alpha) = \sup \{ h_{KS}(\mu), \mu \in \mathcal{M}_{1/2}, \int q d\mu = \alpha \}$. For any open interval $I \subset \mathbb{R}$, we have

$$\liminf_{T \rightarrow +\infty} \frac{\log L_{1/2} \{ \rho \in S^*M, \langle q \rangle_T(\rho) \in I \}}{T} \geq \sup \{ H(\alpha), \alpha \in I \} - (d - 1).$$

The pressure and entropy functions also appear when counting closed geodesics γ with a given q -average:

$$\text{Pr}(q) = \lim_{t \rightarrow +\infty} \frac{1}{t} \log \sum_{\gamma, \ell_\gamma \leq t} e^{\int q d\gamma},$$

and as consequence

$$\limsup_{t \rightarrow +\infty} \frac{1}{t} \log \# \left\{ \gamma, \ell_\gamma \leq t, \int q d\mu_\gamma \in I \right\} \leq \sup \{ H(\alpha), \alpha \in I \},$$

(for a closed interval I),

$$\liminf_{t \rightarrow +\infty} \frac{1}{t} \log \# \left\{ \gamma, \ell_\gamma \leq t, \int q d\mu_\gamma \in I \right\} \geq \sup \{ H(\alpha), \alpha \in I \},$$

(for an open interval I). See [Ki2].

In negative variable curvature, we will also need the following variant of Theorem 3.2.

Theorem 3.3. *Let M be a compact manifold of negative sectional curvature. Let q be a smooth function on S^*M . Let ϕ be a smooth positive function. Denote $L_{1/2}$ the Liouville measure on S^*M . For $\rho \in S^*M$ and $t \in \mathbb{R}$, define $\mathcal{T}_\rho(t)$ by $\int_0^{\mathcal{T}_\rho(t)} \phi(G^s \rho) ds = (d - 1)t$. For $t > 0$, define the function*

$$\langle q \rangle_{\mathcal{T}(-t/2), \mathcal{T}(t/2)} = \frac{1}{\mathcal{T}_\rho(t/2) - \mathcal{T}_\rho(-t/2)} \int_{\mathcal{T}_\rho(-t/2)}^{\mathcal{T}_\rho(t/2)} q \circ G^s(\rho) ds$$

on S^*M .

Then, for any closed interval $I \subset \mathbb{R}$, we have

$$\begin{aligned} \limsup_{t \rightarrow +\infty} \frac{\log L_{1/2} \{ \rho \in S^*M, \langle q \rangle_{\mathcal{T}(-t/2), \mathcal{T}(t/2)}(\rho) \in I \}}{t} \\ \leq (d - 1) \sup \left\{ \frac{h_{KS}(\mu)}{\int \phi d\mu} - \frac{\int \varphi d\mu}{\int \phi d\mu}, \mu \in \mathcal{M}_{1/2}, \int q d\mu \in I \right\}. \end{aligned}$$

For any open interval $I \subset \mathbb{R}$, we have

$$\begin{aligned} \liminf_{t \rightarrow +\infty} \frac{\log L_{1/2} \{ \rho \in S^*M, \langle q \rangle_{\mathcal{T}(-t/2), \mathcal{T}(t/2)}(\rho) \in I \}}{t} \\ \geq (d - 1) \sup \left\{ \frac{h_{KS}(\mu)}{\int \phi d\mu} - \frac{\int \varphi d\mu}{\int \phi d\mu}, \mu \in \mathcal{M}_{1/2}, \int q d\mu \in I \right\}. \end{aligned}$$

Proof. Define a flow \bar{G} on S^*M that has the same trajectories as G but different speed: $\bar{G}^t(\rho) = G^{\mathcal{T}_\rho(t)}(\rho)$. For the new flow, the infinitesimal unstable Jacobian is equal to $(d - 1) \frac{\int \varphi d\mu}{\int \phi d\mu}$. If μ is an invariant probability measure of G , then $d\bar{\mu} = \frac{\phi d\mu}{\int \phi d\mu}$ is an invariant probability measure of \bar{G} . Besides, their entropies are related by the Abramov formula:

$$h_{KS}(\bar{\mu}) = (d - 1) \frac{h_{KS}(\mu)}{\int \phi d\mu},$$

where the entropies of $\bar{\mu}$ and μ are computed with respect to \bar{G} and G respectively.

The theorem is then again an application of Theorem 3.4 in [Ki1] for the Anosov flow \bar{G} . □

4 Averaging

We are now ready to start the proof of Theorem 1.4. We will work in dimension d and constant negative -1 . The changes to make in order to get Theorem 1.7 are indicated in section 7.

The following proposition is proved in [Sj, §2]. It is based on the fact that conjugating $P + i\hbar Q(z)$ by a pseudodifferential operator does not change the spectrum, nor the principal symbol of P . But the “subprincipal symbol” q_z is modified by a coboundary (a function that is cohomologous to zero).

PROPOSITION 1. *Let $T > 0$, there exists an invertible selfadjoint pseudodifferential operator $A_T \in \Psi DO^0$ such that*

$$A_T^{-1} (P + i\hbar Q(z)) A_T = P + i\hbar \text{Op}_\hbar(q^T(z)) + \hbar^2 R_T(z)$$

for $z \in \Omega$; with $R_T \in \Psi DO^0$ depending holomorphically on $z \in \Omega$, and with $q^T(z) \in S^1$ depending holomorphically on $z \in \Omega$, equal to $\langle q \rangle_T - q + q_z$ in a neighbourhood of $p_o^{-1}(1/2)$.

The definition of our symbol classes S^m and operator classes ΨDO^m is given in section 10.

We recall that the operator A_T constructed by Sjöstrand is $A = \text{Op}_\hbar(e^{g_T})$, where

$$g_T = \frac{1}{2} \int_0^{T/2} \left(\frac{2s}{T} - 1 \right) q \circ G^s ds + \frac{1}{2} \int_{-T/2}^0 \left(\frac{2s}{T} + 1 \right) q \circ G^s ds$$

on $p_o^{-1}(1/2)$. The function g_T solves $\{p_o, g_T\} = q - \langle q \rangle_T$. Exactly the same proof yields

PROPOSITION 2. *Assume M has constant curvature -1 . Let $\epsilon > 0$ and $T = (1 - 4\epsilon)|\log \hbar|$. Define $\delta = \frac{1-\epsilon}{2}$. There exists an invertible selfadjoint pseudodifferential operator $A_T \in \Psi DO_\delta^0$ such that*

$$A_T^{-1}(P + i\hbar Q(z))A_T = P + i\hbar \text{Op}_\hbar(q^T(z)) + \hbar^2 R_T(z)$$

for $z \in \Omega$; with $R_T \in \hbar^{-2\delta} \Psi DO_\delta^0$ depending holomorphically on z , and with $q^T(z) \in S_\delta^1$ depending holomorphically on z , equal to $\langle q \rangle_T - q + q_z$ in a neighbourhood of $p_o^{-1}(1/2)$.

In what follows, we will restrict our attention to a region where $|z - \frac{1}{2}| = \mathcal{O}(\hbar)$. As a consequence, we can write

$$P + i\hbar \text{Op}_\hbar(q^T(z)) + \hbar^2 R_T(z) = P + i\hbar \text{Op}_\hbar(q^T) + \hbar \tilde{R}_T(z) \tag{4.1}$$

with $q^T = q^T(1/2) = \langle q \rangle_T$ in a neighbourhood of $p_o^{-1}(1/2)$, and $\tilde{R}_T(z)$ is a pseudodifferential operator depending holomorphically on $z \in \Omega$, tending to zero when $\hbar \rightarrow 0$ and $|z - \frac{1}{2}| = \mathcal{O}(\hbar)$. More precisely,

$$\tilde{R}_T(z) = \left(z - \frac{1}{2}\right) Q'(z) + \hbar R_T(z),$$

$R_T \in \hbar^{-2\delta} \Psi DO_\delta^0$ depending holomorphically on z , and $Q'(z) \in \Psi DO_\delta^1$ depending holomorphically on z .

In the following sections, we let z vary in a disc of radius $\mathcal{O}(\hbar)$ around $1/2$. We will write $2z = 1 + \zeta$, $\zeta = \mathcal{O}(\hbar)$. We consider the operator (4.1), that we write

$$\mathcal{P}_T = \mathcal{P}_T(z) = P + i\hbar Q_T + \hbar \tilde{R}_T(z), \quad Q_T = \text{Op}_\hbar(q^T). \tag{4.2}$$

Note that we have

$$\{p_o, q^T\} = \mathcal{O}\left(\frac{1}{T}\right)$$

in a neighbourhood of $p_o^{-1}(1/2)$. By Proposition 8, this implies

$$\|[P, Q_T]u\| \leq C \left(\frac{\hbar}{T} + \mathcal{O}(\hbar^{2-2\delta}) \right) \|u\| + \mathcal{O}(\hbar) \left\| \left(P - \frac{1}{2}\right) u \right\|. \tag{4.3}$$

5 Perturbations with Controlled Trace Norm

We now want to make a small perturbation $\tilde{\mathcal{P}}$ of \mathcal{P} with a good control over the resolvent $(\tilde{\mathcal{P}}(z) - z)^{-1}$, and over the trace class norm $\|\cdot\|_1$ of $\tilde{\mathcal{P}} - \mathcal{P}$. The general

idea is the following: we want to count the eigenvalues of \mathcal{P} in a given open set Ω . Suppose that, by perturbing \mathcal{P} into $\tilde{\mathcal{P}}$, we can make sure that $\tilde{\mathcal{P}}$ has no eigenvalue in $\tilde{\Omega} \supset \Omega$; and suppose that we control the trace class norm of the perturbation $\tilde{\mathcal{P}} - \mathcal{P}$. Then this information should suffice to give an upper bound on the number of eigenvalues of the original operator \mathcal{P} , inside Ω . The full argument will be given in section 6.

We construct a pseudodifferential operator $\tilde{Q}_T = \text{Op}_h(\tilde{q}^T) \in \Psi DO_\delta^1$ such that $\tilde{q}^T \leq q^T$ on $p_o^{-1}(1/2)$ and $\{p_o, \tilde{q}^T\} = \mathcal{O}(1/T)$. In addition, we fix some $\epsilon > 0$ and introduce an arbitrarily small $\theta > 0$, and we want $q^T = \tilde{q}^T$ on $p_o^{-1}(\frac{1}{2} - \epsilon, \frac{1}{2} + \epsilon) \cap \{q^T \leq \alpha - 3\theta\}$, and $\tilde{q}^T \leq \alpha - 2\theta$ everywhere on $p_o^{-1}(\frac{1}{2} - \epsilon, \frac{1}{2} + \epsilon)$. For instance we can take $\tilde{q}^T = a(q^T)$ on $p_o^{-1}(\frac{1}{2} - \epsilon, \frac{1}{2} + \epsilon)$ where a is real and smooth, $a(E) \leq E$, $|a'| \leq 1$; $a(E) = E$ if $E \leq \alpha - 3\theta$, and $a \leq \alpha - 2\theta$ everywhere.

REMARK 5.1. At this stage it is convenient to choose a positive quantization scheme Op_h , in order to have $\text{Op}_h(q^T) \geq \text{Op}_h(\tilde{q}^T)$.

Let $0 \leq f \in \mathcal{S}(\mathbb{R})$, with $\hat{f} \in C_0^\infty$, where $\hat{f}(t) = \int e^{itE} f(E) dE$ is the Fourier transform. Put

$$\tilde{\mathcal{P}} = P + i\hbar\hat{Q}_T + \hbar\tilde{R}_T(z),$$

with

$$\hat{Q}_T = Q_T + f\left(\frac{2P-1}{\hbar}\right)(\tilde{Q}_T - Q_T)f\left(\frac{2P-1}{\hbar}\right).$$

The following proposition is proved in [Sj] for fixed T (and $\delta = 0$, that is, with standard symbol classes). One can follow the proof of [Sj] line by line and check that it is still valid for $T = (1 - 4\epsilon)|\log \hbar|$, $\epsilon > 0$ very small.

PROPOSITION 3. Let $P = -\hbar^2 \frac{\Delta}{2}$. Let $Q = Q(z) \in \Psi DO^1$ have principal symbol $q(z)$ depending holomorphically on $z \in \Omega$, and be formally self-adjoint when z is real. Let

$$\mathcal{P}_T = P + i\hbar Q_T + \hbar\tilde{R}_T(z), \quad Q_T = Q_T\left(\frac{1}{2}\right), \quad z = \frac{1+\zeta}{2}, \quad \zeta = \mathcal{O}(\hbar),$$

be the operator defined in (4.2), with $Q_T = \text{Op}_h(q^T) \in \Psi DO_\delta^1$, and $\tilde{R}_T(z) \in \hbar^{1-2\delta}\Psi DO_\delta^0 + (z - \frac{1}{2})\Psi DO_\delta^1$. Let $\tilde{Q}_T = \text{Op}_h(\tilde{q}^T) \in \Psi DO_\delta^1$, with $\tilde{q}^T = a(q^T)$ on $p_o^{-1}(\frac{1}{2} - \epsilon, \frac{1}{2} + \epsilon)$, where a is real and smooth, $a(E) \leq E$, $|a'| \leq 1$; $a(E) = E$ if $E \leq \alpha - 3\theta$, and $a \leq \alpha - 2\theta$.

Put

$$\tilde{\mathcal{P}}_T = P + i\hbar\hat{Q}_T + \hbar\tilde{R}_T(z),$$

with

$$\hat{Q}_T = Q_T + f\left(\frac{2P-1}{\hbar}\right)(\tilde{Q}_T - Q_T)f\left(\frac{2P-1}{\hbar}\right).$$

Then

$$\|\tilde{\mathcal{P}}_T - \mathcal{P}_T\| \leq \hbar\left(\|f\|_\infty^2 \sup_{p_o^{-1}(1/2)} (q^T - \tilde{q}^T) + \mathcal{O}(\hbar^{1-2\delta})\right)$$

and for any integer N ,

$$\begin{aligned} \|\tilde{\mathcal{P}}_T - \mathcal{P}_T\|_1 &\leq C_d \hbar^{2-d} \left[\hat{f}^2(0) \int_{p_o^{-1}(1/2)} (q^T - \tilde{q}^T) L_{1/2}(d\rho) \right. \\ &\quad \left. + \sum_{k=1}^{N-1} \hbar^k |D_t^{2k} \hat{f}^2(0)| \int_{p_o^{-1}(1/2)} |D_\rho^{2k} (q^T - \tilde{q}^T)| L_{1/2}(d\rho) + \mathcal{O}(\hbar^{N(1-2\delta)}) \right], \end{aligned} \quad (5.1)$$

where D_t^{2k} and D_ρ^{2k} are differential operators of degree $\leq 2k$, respectively on \mathbb{R} and T^*M .

If we restrict z by assuming that for some continuous function $F(E) > 0$, defined on some bounded interval J containing 0, we have

$$\frac{\Im m(\zeta)}{2\hbar} - q^T + f\left(\frac{\Re e(\zeta)}{\hbar}\right)^2 (q^T - \tilde{q}^T) \geq F\left(\frac{\Re e(\zeta)}{\hbar}\right),$$

on $p_o^{-1}(1/2)$, $\Re e(\zeta)/\hbar \in J$, then for \hbar small enough, $(z - \tilde{\mathcal{P}}_T)^{-1}$ exists, and we have

$$\left\| \left(\frac{1}{\hbar} (z - \tilde{\mathcal{P}}_T) \right)^{-1} \right\| \leq \frac{2 + 12 \sup_{p_o^{-1}(1/2)} (q^T - \tilde{q}^T) \|f'\|_\infty \|f\|_\infty}{F(\Re e(\zeta)/\hbar)}.$$

The proof is identical to the proof in [Sj]. In Appendix 11 we will give some details, for the reader's convenience. Taking $F(E) = \min\{\theta/2, 3\theta f^2(E)/2\}$, we find

COROLLARY 2. *Define*

$$\tilde{\Omega}_\hbar = \left\{ \frac{1}{2} - 2c\hbar \leq \Re e(z) \leq \frac{1}{2} + 2c\hbar \right\} \cap \left\{ (\alpha - \theta)\hbar \leq \Im m(z) \leq 4\|q\|_\infty \hbar \right\} \subset \mathbb{C}.$$

For $z \in \tilde{\Omega}_\hbar$, the operator $z - \tilde{\mathcal{P}}_T$ is invertible, and

$$\|(\tilde{\mathcal{P}}_T - z)^{-1}\| \leq \frac{C_{f,q}}{\theta\hbar}.$$

As a direct consequence of (5.1), we obtain

COROLLARY 3. *For \hbar small enough, we have*

$$\begin{aligned} \|\tilde{\mathcal{P}}_T - \mathcal{P}_T\|_1 &\leq C_{d,f,q} \hbar^{2-d} L_{1/2} \left(\{\tilde{q}^T \neq q^T\} \cap \{p_o^{-1}(\tfrac{1}{2})\} \right) \\ &\leq C_{d,f,q} \hbar^{2-d} L_{1/2} \left(\{q^T \geq \alpha - 3\theta\} \cap \{p_o^{-1}(\tfrac{1}{2})\} \right) \\ &\leq C_{d,f,q} \hbar^{2-d} e^{T[H(\alpha-3\theta)-(d-1)+\epsilon]} \\ &\leq C_{d,f,q} \hbar^{2-d} \hbar^{[(d-1)-H(\alpha-3\theta)-\epsilon](1-4\epsilon)} \end{aligned}$$

for \hbar small enough.

We recall that $\theta > 0$ and $\epsilon > 0$ can be chosen arbitrarily small. The two first inequalities come from (5.1), which is Sjöstrand's estimate pushed to time $T = (1 - 4\epsilon)|\log \hbar|$. The third inequality is a consequence of Kifer's theorem 3.2.

6 Jensen's Inequality

We can now implement the strategy outlined in the previous section. The spectrum of \mathcal{P} will be expressed as the set of zeroes of a certain determinant, which is a holomorphic function. We will use the Jensen inequality to count these zeroes.

We already defined

$$\tilde{\Omega}_h = \left\{ \frac{1}{2} - 2c\hbar \leq \Re e(z) \leq \frac{1}{2} + 2c\hbar \right\} \cap \left\{ (\alpha - \theta)\hbar \leq \Im m(z) \leq 4\|q\|_\infty \hbar \right\} \subset \mathbb{C}.$$

To finish the proof of Theorem 1.4, we also introduce the set

$$\Omega_h = \left\{ \frac{1}{2} - c\hbar \leq \Re e(z) \leq \frac{1}{2} + c\hbar \right\} \cap \left\{ \alpha\hbar \leq \Im m(z) \leq 3\|q\|_\infty \hbar \right\} \subset \tilde{\Omega}_h.$$

For $z \in \tilde{\Omega}_h$, we can write

$$\mathcal{P}_T - z = (\tilde{\mathcal{P}}_T - z)(1 + K(z))$$

where $K(z)$ is the trace class operator $(\tilde{\mathcal{P}}_T - z)^{-1}(\mathcal{P}_T - \tilde{\mathcal{P}}_T)$.

We can bound the number of eigenvalues of \mathcal{P}_T in Ω_h by the number of zeros of the holomorphic function $g(z) = \det(1 + K(z))$ in Ω_h . Let us call $N(g, \Omega_h)$ this number of zeros. Introduce $z_0 = \frac{1}{2} + 2i\hbar\|q\|_\infty$. By the Jensen inequality [Ru],

$$N(g, \Omega_h) \leq C \left(\log \|g\|_{\infty, \tilde{\Omega}_h} - \log |g(z_0)| \right), \tag{6.1}$$

where the constant C does not depend on \hbar (because the rectangles $\tilde{\Omega}_h$ and Ω_h can be transported, by translations and homotheties, to the fixed rectangles $\tilde{\Omega}_1$ and Ω_1).

On the one hand, for all $z \in \tilde{\Omega}_h$,

$$\begin{aligned} |\det(1 + K(z))| &\leq \exp \|K(z)\|_1 \\ &\leq \exp \left(\|(\tilde{\mathcal{P}}_T - z)^{-1}\| \|\mathcal{P}_T - \tilde{\mathcal{P}}_T\|_1 \right) \\ &\leq \exp \left(\frac{C_{d,f,q}}{\theta\hbar} \hbar^{2-d} \hbar^{[(d-1)-H(\alpha-3\theta)-\epsilon](1-4\epsilon)} \right) \\ &\leq \exp \left(C_{f,q,\theta,d} \hbar^{1-d} \hbar^{[(d-1)-H(\alpha-3\theta)-\epsilon](1-4\epsilon)} \right). \end{aligned}$$

On the other hand, we also know that $\|(1 + K(z_0))^{-1}\| \leq C\hbar^{-1}$: since z_0 has ‘large’ imaginary part, $\mathcal{P}_T - z_0$ is invertible, and it is easy to get a bound $\|(\mathcal{P}_T - z_0)^{-1}\| = \mathcal{O}(\hbar^{-1})$. We use the same calculation as in [Sj] and get

$$\begin{aligned} |\det(1 + K(z_0))^{-1}| &= |\det(1 - K(z_0)(1 + K(z_0))^{-1})| \\ &\leq \exp \|K(z_0)(1 + K(z_0))^{-1}\|_1 \\ &\leq \exp \|K(z_0)\|_1 \|(1 + K(z_0))^{-1}\| \\ &\leq \exp \left(\tilde{C}_{f,q,\theta,d} \hbar^{1-d} \hbar^{[(d-1)-H(\alpha-3\theta)-\epsilon](1-4\epsilon)} \right) \end{aligned}$$

so that

$$|\det(1 + K(z_0))| \geq \exp \left(-\tilde{C}_{f,q,\theta,d} \hbar^{1-d} \hbar^{[(d-1)-H(\alpha-3\theta)-\epsilon](1-4\epsilon)} \right).$$

This, combined to (6.1), yields

$$|N(g, \Omega_h)| \leq C \hbar^{1-d} \hbar^{[(d-1)-H(\alpha-2\theta)-\epsilon](1-4\epsilon)}$$

Since θ and $\epsilon > 0$ are arbitrary, we have proved Theorem 1.4.

7 The Case of Variable Negative Curvature

We now indicate how to modify the arguments in order to treat the case of variable curvature (Theorem 1.7).

We should modify Proposition 2 as follows. Let $T = (1 - 4\epsilon)|\log \hbar|$ as before. Fix ϕ a smooth function such that $\phi \geq \varphi$. Define $\mathcal{T}_\rho(T/2)$, $\mathcal{T}_\rho(-T/2)$ as in Theorem 3.3, in a neighbourhood of $p_o^{-1}(1/2)$. We have to choose ϕ smooth because we want \mathcal{T}_ρ to depend smoothly on ρ .

In dimension $d = 2$, we have

$$\|d_\rho G^t\| \leq e^{\frac{T}{2}} \leq \hbar^{-\frac{1-4\epsilon}{2}}$$

for $\mathcal{T}_\rho(-T/2) \leq t \leq \mathcal{T}_\rho(T/2)$. This does not hold for $d > 2$ and this is why our proof only works for surfaces. More generally, $\|d_\rho^k G^t\| \leq e^{\frac{kT}{2}} \leq \hbar^{-\frac{(1-4\epsilon)k}{2}}$ for derivatives of higher order.

The derivatives of $\rho \mapsto \mathcal{T}_\rho(\frac{T}{2})$ can be computed by the implicit function theorem, and it follows that $\mathcal{T}_\rho(\frac{T}{2}) \in S_\delta^0$ with $\delta = \frac{1-\epsilon}{2}$. In Proposition 2, we now define $A_T = \text{Op}_\hbar(e^{g_T})$ where

$$g_T(\rho) = \frac{1}{2} \int_0^{\mathcal{T}_\rho(T/2)} \left(\frac{s}{\mathcal{T}_\rho(T/2)} - 1 \right) q \circ G^s ds + \frac{1}{2} \int_{-\mathcal{T}_\rho(T/2)}^0 \left(\frac{s}{\mathcal{T}_\rho(T/2)} + 1 \right) q \circ G^s ds$$

on $p_o^{-1}(1/2)$. In the last sentence of Proposition 2, we replace $\langle q \rangle_T$ by $\langle q \rangle_{\mathcal{T}(-T/2), \mathcal{T}(T/2)} + r_T$ where $r_T = q - \{p_o, g_T\} - \langle q \rangle_{\mathcal{T}(-T/2), \mathcal{T}(T/2)}$ satisfies $r_T \in |\log \hbar|^{-1} S_\delta^0$ and $\{p_o, r_T\} \in |\log \hbar|^{-1} S_\delta^0$. For $d = 2$, all the operators A_T , R_T , etc., stay in the same class as stated in Proposition 2.

Starting from here, the proof of Theorem 1.7 goes exactly along the same lines as before. We find

$$\limsup_{\hbar \rightarrow 0} \frac{\log \#\{z \in \Sigma_{1/2}, \frac{\Im m(z)}{\hbar} \geq \alpha\}}{|\log \hbar|} \leq \tilde{H}(\alpha),$$

where $\tilde{H}(\alpha) = (d - 1) \sup \left\{ \frac{h_{KS}(\mu)}{\int \phi d\mu} - \frac{\int \varphi d\mu}{\int \phi d\mu} + 1, \mu \in \mathcal{M}_{1/2}, \int q d\mu = \alpha \right\}$. Letting ϕ converge to φ uniformly, we obtain Theorem 1.7.

8 About Question 3 (Q3)

In this section, we consider a particular case of the spectral problem (1.5) in which the trace formula is exact. We then try to investigate (Q3) on this example.

Let M be a compact hyperbolic surface: M can be written as $M = \Gamma \backslash \mathbb{H}$, where \mathbb{H} is the hyperbolic disc and Γ is a discrete subgroup of the group of hyperbolic isometries. Let $[\omega] \in H^1(M, \mathbb{C})$ be represented by the harmonic complex valued 1-form ω . Introduce the twisted laplacian

$$\Delta_\omega f = \Delta f - 2\langle \omega, df \rangle + \|\omega\|_x^2 f.$$

Studying the large eigenvalues of Δ_ω amounts to studying a fixed spectral window for the semiclassical twisted laplacian

$$-\hbar^2 \frac{\Delta \omega}{2} = -\hbar^2 \frac{\Delta}{2} + \hbar^2 \langle \omega, d. \rangle - \hbar^2 \frac{\|\omega\|_x^2}{2}, \quad \hbar \rightarrow 0.$$

The ‘‘usual’’ selfadjoint case is when ω has coefficient in $i\mathbb{R}$. We shall instead be interested in the case when ω has coefficients in \mathbb{R} . The operator falls exactly

into the case studied in section 1.2, with $q(x, \xi) = \langle \omega_x, \xi \rangle$. The geodesic flow is ergodic, and Sjöstrand's result tells us that "most" eigenvalues of $-\hbar^2 \frac{\Delta_\omega}{2}$ such that $\Re e(z) \in [\frac{1}{2} - C\hbar, \frac{1}{2} + C\hbar]$ have imaginary part $\Im m(z) = o(\hbar)$. Equivalently, "most" eigenvalues of $-\Delta_\omega$ such that $\Re e(z) \in [\lambda - C\sqrt{\lambda}, \lambda + C\sqrt{\lambda}]$ have imaginary part $\Im m(z) = o(\sqrt{\lambda})$.

We rephrase (Q3) as

(Q3') If $\omega \neq 0$, is it possible to have $\Im m(z)/\hbar \rightarrow 0$ as $\hbar \rightarrow 0$ and $\Re e(z) \in [\frac{1}{2} - C\hbar, \frac{1}{2} + C\hbar]$, $z \in Sp(-\hbar^2 \frac{\Delta_\omega}{2})$?

CONJECTURE. I conjecture the opposite: if $\omega \neq 0$, then there is a sequence $\hbar_n \rightarrow 0$, $z_n \in Sp(-\hbar_n^2 \frac{\Delta_\omega}{2})$ with $\Re e(z_n) \in [\frac{1}{2} - C\hbar_n, \frac{1}{2} + C\hbar_n]$ and $\frac{\Im m(z_n)}{\hbar_n} \not\rightarrow 0$.

As is usual in hyperbolic spectral theory, we introduce the spectral parameter r : if λ_j is an eigenvalue of $-\Delta_\omega$, we denote $\lambda_j = \frac{1}{4} + r_j^2$. Yet another way to phrase (Q3) is to ask whether $\Im m(r_j) \rightarrow 0$ as $\Re e(r_j) \rightarrow \infty$. Sjöstrand's results say that $\Im m(r_j)$ is bounded and that $\Im m(r_j) \rightarrow 0$ for a subsequence of density one. But I believe that it is impossible to have $\Im m(r_j) \rightarrow 0$ for the whole sequence, unless $\omega = 0$.

Recall the Selberg trace formula [Se] valid for $\omega \in H^1(M, i\mathbb{R})$,

$$\sum_{\lambda_j} \hat{f}(r_j) = \frac{Area(M)}{4\pi} \int_{-\infty}^{+\infty} \hat{f}(r) r \tanh(\pi r) dr + \sum_{\gamma} \frac{e^{\int_{\gamma} \omega} \ell_{\gamma_o}}{\sinh \frac{\ell_{\gamma}}{2}} f(\ell_{\gamma}), \tag{8.1}$$

for any function f on \mathbb{R} , even, smooth enough, and decaying faster than any exponential. On the right-hand side, the sum runs over the set of closed geodesics (equivalently, the set of conjugacy classes in Γ). If γ is a periodic geodesic, we denote $\ell_{\gamma} > 0$ its length, or period; and ℓ_{γ_o} is its shortest period.

Recall that the Fourier transform is defined by $\hat{f}(r) = \int e^{iru} f(u) du$.

PROPOSITION 4. *The trace formula holds, under the same assumptions on f , if $\omega \in H^1(M, \mathbb{R})$.*

Proof. Take $\omega \in H^1(M, \mathbb{R})$. We consider the operator $\Delta_{z\omega}$ for $z \in \mathbb{C}$. The argument of section 2 shows that this operator has discrete spectrum (eigenvalues).

The right-hand side of the trace formula reads

$$\frac{Area(M)}{4\pi} \int_{-\infty}^{+\infty} \hat{f}(r) r \tanh(\pi r) dr + \sum_{\gamma} \frac{e^{z \int_{\gamma} \omega} \ell_{\gamma_o}}{\sinh \ell_{\gamma}/2} f(\ell_{\gamma}) \tag{8.2}$$

and clearly is an entire function of z .

The left-hand side is $\sum_j \hat{f}(r_j(z))$. To check that it is an entire function of z , we first note that $\hat{f}(r)$, being an even entire function, can be written as $g(\frac{1}{4} + r^2)$ where g is entire. Thus $\hat{f}(r_j(z)) = g(\lambda_j(z))$, where the $\lambda_j(z)$ are the eigenvalues of $-\Delta_{z\omega}$.

If z is restricted to a bounded subset Ω of \mathbb{C} , we note that the $\Im m(r_j(z))$ are uniformly bounded. To that end, we write $-\lambda_j = \frac{1}{4} + r_j^2$ with $r_n = x + iy$. The eigenvalue equation

$$-\Delta_{z\omega} f = \lambda_j f$$

with f normalized in L^2 implies both equations

$$\frac{1}{4} + x^2 - y^2 = \int |\nabla f|^2 + 2\beta i \int \langle \omega, df \rangle \bar{f} + (\beta^2 - \alpha^2) \int \|\omega\|_x^2 |f|^2$$

and

$$2xy = -2\alpha i \int \langle \omega, df \rangle \bar{f} - 2\alpha\beta \int \|\omega\|_x^2 |f|^2$$

if we decompose $z = \alpha + i\beta \in \mathbb{R} + i\mathbb{R}$ and $r_j = x + iy \in \mathbb{R} + i\mathbb{R}$. If α and β stay bounded, it also follows that y must stay bounded.

Besides

$$\lambda^{-2\#} \{n, 0 \leq \Re e(r_n(z)) < \lambda\}$$

is bounded uniformly for $\lambda > 1$ and z staying in a compact set of \mathbb{C} (the arguments of [Sj, §4], or of our sections 4, 5 and 6 are locally uniform in z). Since \hat{f} is rapidly decreasing in each horizontal strip, it follows that the sum $\sum_j \hat{f}(r_j(z))$ is the uniform limit of the partial sums $\sum_{|\Re e(r_j(z))| < \lambda} \hat{f}(r_j(z))$. But for a given λ , this is a holomorphic function of z (in the open set $\{z, \Re e(r_j(z)) \neq \lambda \text{ for all } j\}$), since $\Delta_{z\omega}$ depends holomorphically on z and $\sum_{|\Re e(r_j(z))| < \lambda} \hat{f}(r_j(z))$ can be defined by holomorphic functional calculus.

This shows that $\sum_j \hat{f}(r_j(z))$ is also an entire function. Both sides of (8.2) coincide for $z \in i\mathbb{R}$ (by the usual trace formula), thus they must coincide for all $z \in \mathbb{C}$. \square

Introduce some parameters $\sigma, R, T > 0$, and take

$$f(u) = \frac{1}{\sqrt{2\pi\sigma}} \left[e^{-\frac{(u-T)^2}{2\sigma^2}} e^{iuR} + e^{-\frac{(u+T)^2}{2\sigma^2}} e^{-iuR} \right]$$

so that

$$\hat{f}(r) = e^{-\frac{\sigma^2}{2}(r-R)^2} e^{-iTr} + e^{-\frac{\sigma^2}{2}(r+R)^2} e^{iTr}.$$

This yields

$$\begin{aligned} & \sum_j e^{-\frac{\sigma^2}{2}(r_j-R)^2} e^{-iTr_j} + e^{-\frac{\sigma^2}{2}(r_j+R)^2} e^{iTr_j} \\ &= \frac{\text{Area}(M)}{4\pi} \int_{-\infty}^{+\infty} r \tanh \pi r \left[e^{-\frac{\sigma^2}{2}(r-R)^2} e^{-iTr} + e^{-\frac{\sigma^2}{2}(r+R)^2} e^{iTr} \right] dr \\ & \quad + \sum_{\gamma} \frac{e^{\int_{\gamma} \omega} \ell_{\gamma_0}}{\sinh \frac{\ell_{\gamma}}{2}} \frac{1}{\sqrt{2\pi\sigma}} \left[e^{-\frac{(\ell_{\gamma}-T)^2}{2\sigma^2}} e^{i\ell_{\gamma}R} + e^{-\frac{(\ell_{\gamma}+T)^2}{2\sigma^2}} e^{-i\ell_{\gamma}R} \right]. \end{aligned} \tag{8.3}$$

We want to bound the right-hand side from below. We hope that this bound will tell us that $\exp(\pm iTr_j)$ cannot be too small on the left-hand side, giving some information on the imaginary part of r_j , for $\Re e(r_j) \sim R$. On the right, the idea is that the main contribution should come from the geodesics with $\ell_{\gamma} \sim T$. We want to choose R so as not to be bothered by the oscillatory terms $e^{\pm i\ell_{\gamma}R}$. For that purpose, R and T will be related in the following manner.

LEMMA 8.1 [PhR, Lem. 3.3],[JP]. *For any $M > 1$, there exists*

$$R \in [M, M \exp(\exp(5T))]$$

such that $\cos(R\ell_{\gamma}) \geq 1/2$ for every closed geodesic γ with $\ell_{\gamma} \leq 5T$.

In the sequel, we take $M = \exp(\exp(cT))$ ($c > 0$ arbitrary) to ensure that T is of order $\log \log R$. We note that this relation between R and T is independent of ω . This will allow us to modify slightly our initial problem by extending it to the case where ω can depend on R (or T). More precisely, we want to consider the case when $\omega = \Theta(R)\omega_o$, where ω_o is fixed and $\Theta(R) \geq 1$ is allowed to go to infinity with R at a reasonable rate.

This means that we consider a slight generalization of (1.5) (the motivation should become clearer in section 8.2).

8.1 A more general problem. We consider a spectral problem of the form

$$(\mathcal{P} - z)u = 0 \tag{8.4}$$

where

$$\mathcal{P} = \mathcal{P}(z) = P + i\hbar\theta(\hbar)Q(z), \quad P = -\frac{\hbar^2\Delta}{2},$$

where $z \in \Omega = e^{i[-s_0, s_0]}[E_{\min}, E_{\max}[$, with $0 < E_{\min} < 1/2 < E_{\max} < +\infty$, $0 < s_0 < \pi/4$. We will assume that $Q(z) \in \Psi DO^1$ depends holomorphically on $z \in \Omega$, and that $\theta(\hbar)$ is some real valued function such that $\theta(\hbar) \geq 1$ and $\hbar\theta(\hbar) \xrightarrow{\hbar \rightarrow 0} 0$. We have in mind $\theta(\hbar) = |\log(\hbar)|$. Finally, we assume that Q is formally self-adjoint for z real. Again, we call Σ the ‘‘spectrum’’ the set of z for which the equation $(\mathcal{P}(z) - z)u = 0$ has a solution.

The results of section 1.2 can be generalized as follows: for any $E_{\min} < E_1 \leq E_2 < E_{\max}$,

$$\#\{z \in \Sigma, E_1 \leq \Re e(z) \leq E_2\} = \frac{1}{(2\pi\hbar)^d} \left[\int_{p_o^{-1}[E_1, E_2]} dx d\xi + \mathcal{O}(\hbar\theta(\hbar)) \right]. \tag{8.5}$$

One can show that $\Im m(z)/\hbar\theta(\hbar)$ has to stay bounded for $z \in \Sigma$. Taking $\theta(\hbar) = |\log(\hbar)|$, $E_1 = E - c\hbar\theta(\hbar)$ and $E_2 = E + c\hbar\theta(\hbar)$, one has

$$q_E^- + o(1) \leq \frac{\Im m(z)}{\hbar\theta(\hbar)} \leq q_E^+ + o(1)$$

for $z \in \Sigma$ such that $E_1 \leq \Re e(z) \leq E_2$. Assuming that the geodesic flow is ergodic on $p_o^{-1}\{E\}$, and taking $\theta(\hbar) = |\log(\hbar)|$, we have for any $\epsilon > 0$, any $c > 0$,

$$\#\left\{z \in \Sigma, E - c\hbar\theta(\hbar) \leq \Re e(z) \leq E + c\hbar\theta(\hbar), \frac{\Im m(z)}{\hbar\theta(\hbar)} \notin [\bar{q}_E - \epsilon, \bar{q}_E + \epsilon]\right\} = \theta(\hbar)o(\hbar^{1-d}). \tag{8.6}$$

REMARK 8.2. The paper [Sj] only treats the case $\theta(\hbar) = 1$. But the method of [Sj, §5] can be adapted in a straightforward way to show the following: consider the spectral problem

$$(\mathcal{P} - z)u = 0 \tag{8.7}$$

where

$$\mathcal{P} = \mathcal{P}(z) = P + i\hbar\theta Q(z), \quad P = -\frac{\hbar^2\Delta}{2},$$

where $z \in \Omega = e^{i[-s_0, s_0]}[E_{\min}, E_{\max}[$, with $0 < E_{\min} < 1/2 < E_{\max} < +\infty$,

$0 < s_0 < \pi/4$. Fix some $\epsilon > 0$. Then

$$\#\{z \in \Sigma, E_1 \leq \Re e(z) \leq E_2\} = \frac{1}{(2\pi\hbar)^d} \left[\int_{p_o^{-1}[E_1, E_2]} dx d\xi + \mathcal{O}(\hbar\theta) \right]. \quad (8.8)$$

uniformly in all $\theta \geq 1$ such that $\theta\hbar \leq \epsilon$, and E_1, E_2 such that $E_{\min} < E_1 - 2\epsilon$, $E_2 + 2\epsilon < E_{\max}$, $|E_2 - E_1| \geq \hbar\theta$.

For (8.6), and if we take $\theta(\hbar) = |\log \hbar|$, the generalization of the proof in [Sj] is without surprise, but requires some rather technical changes; we will not prove it here, but still feel allowed to ask about (Q3) in this generalized setting. We note that to extend (8.6) to more general $\theta(\hbar)$, some analyticity assumptions would be required, exactly as in [HiSV].

We will focus our attention on the operator

$$-\hbar^2 \frac{\Delta_{\theta(\hbar)\omega}}{2} = -\hbar^2 \frac{\Delta}{2} + \hbar^2 \theta(\hbar) \langle \omega, d. \rangle - \hbar^2 \theta(\hbar)^2 \frac{\|\omega\|_x^2}{2}, \quad \hbar \rightarrow 0,$$

when ω has coefficients in \mathbb{R} .

We generalize (Q3) to

(Q3'') If $\omega \neq 0$, prove that there exists a sequence $\hbar_n \rightarrow 0$, and $z_n \in Sp(-\hbar_n^2 \frac{\Delta_{\theta(\hbar_n)\omega}}{2})$ with $\Re e(z_n) \in [\frac{1}{2} - C\hbar_n\theta(\hbar_n), \frac{1}{2} + C\hbar_n\theta(\hbar_n)]$, such that $\frac{\Im m(z_n)}{\hbar_n\theta(\hbar_n)} \not\rightarrow 0$.

8.2 Heuristic discussion of the parameters σ, R, T . We start again from the trace formula (8.3), considering the case where $\omega = \Theta(R)\omega_o$, $\Theta(R) = \theta(R^{-1}) \geq 1$. Here R^{-1} is going to play the role of the small parameter \hbar . The form ω_o is fixed, and we normalize it to have stable norm $\|\omega_o\|_s = 1$.

$$\|\omega\|_s = \sup \left\{ \int_{S^*M} \omega d\mu, \mu \in \mathcal{M}_{1/2} \right\} \quad (8.9)$$

$$= \sup_{\gamma} \frac{\int_{\gamma} \omega}{\ell_{\gamma}}. \quad (8.10)$$

The first line can be considered as a definition of the stable norm (valid for a general compact Riemannian manifold M), whereas the second line holds on negatively curved manifolds because of the density of the closed geodesics. Using the definition $\Pr(\omega) = \sup \{h_{KS}(\mu) + \int_{S^*M} \omega d\mu, \mu \in \mathcal{M}_{1/2}\}$, it is not difficult to show that

$$\lim_{t \rightarrow \infty} \Pr(t\omega) - t\|\omega\|_s = \sup \left\{ h_{KS}(\mu), \mu \in \mathcal{M}_{1/2}, \int_{S^*M} \omega d\mu = \|\omega\|_s \right\}.$$

On a surface, the right-hand side vanishes [A2]. Besides, for any T one can find a closed geodesic γ with $\ell_{\gamma} \in [T - 1, T]$, and such that

$$\frac{\int_{\gamma} \omega_o}{\ell_{\gamma}} \geq \|\omega_o\|_s (1 + o_T(1)) = (1 + o_T(1)), \quad (8.11)$$

where $o_T(1)$ goes to 0 as T approaches $+\infty$. Simply recall that for any $0 < \delta < 1$,

$$\lim_{T \rightarrow \infty} \frac{\log \#\{\gamma, \ell_{\gamma} \in [T - 1, T], \int_{\gamma} \omega_o / \ell_{\gamma} \geq (1 - \delta)\}}{T} = H(1 - \delta) > 0, \quad (8.12)$$

where

$$H(\alpha) = \sup \left\{ h_{KS}(\mu), \mu \in \mathcal{M}_{1/2}, \int_{S^*M} \omega_o d\mu = \alpha \right\}.$$

The function H is continuous, concave on $[-1, 1]$, real-analytic on $] - 1, 1[$ ([BL]). And again, $H(-1) = H(1) = 0$ on a compact surface [A2].

In (8.3), we have not said yet how σ will depend on R and T . For the moment, let us decide *a priori* that σ should be such that the term

$$\frac{Area(M)}{4\pi} \int_{-\infty}^{+\infty} r \tanh \pi r \left[e^{-\frac{\sigma^2}{2}(r-R)^2} e^{-iTr} + e^{-\frac{\sigma^2}{2}(r+R)^2} e^{iTr} \right] dr$$

is negligible compared to the sum \sum_{γ} . Remember that R and T are chosen so as to satisfy Lemma 8.1. Then, fixing $1 > \delta > 0$, the right-hand side of (8.3) should be bounded from below by

$$\sigma^{-1} e^{TH(1-\delta)-T/2} e^{\Theta(R)(1-\delta)T} e^{-1/2\sigma^2}. \tag{8.13}$$

We want to use the fact that this grows quite fast with T . On the other hand, looking at the left-hand side of (8.3), we cannot hope to do better than to bound it from above by

$$\#\{j, |\Re(r_j) - R| \leq \sigma^{-1}\} e^{\frac{\sigma^2}{2} \sup_j \Im m(r_j)^2} e^{T \sup_j |\Im m(r_j)|}, \tag{8.14}$$

where each time the \sup_j should be restricted to the indices j such that $|\Re(r_j) - R| \leq \sigma^{-1}$.

This heuristic argument would give an inequality

$$\begin{aligned} \#\{j, |\Re(r_j) - R| \leq \sigma^{-1}\} e^{\frac{\sigma^2}{2} \sup_j \Im m(r_j)^2} e^{T \sup_j |\Im m(r_j)|} \\ \geq \sigma^{-1} e^{TH(1-\delta)-T/2} e^{\Theta(R)(1-\delta)T} e^{-\frac{1}{2\sigma^2}}, \end{aligned} \tag{8.15}$$

obtained by comparing the lower bound (8.13) and the upper bound (8.14). Again, the hope is to compare the powers of e^T on both sides to prove that $\sup_j |\Im m(r_j)|$ cannot be arbitrarily small.

Consider the case $\Theta(R) = 1$, which is the case we were originally interested in. If we take σ to be a constant, then by Weyl's law we have $\#\{j, |\Re(r_j) - R| \leq \sigma^{-1}\} \sim R \geq \exp(\exp(cT))$. In this case (8.15) cannot bring any useful information. On the other hand, if we want to choose σ such that $\#\{j, |\Re(r_j) - R| \leq \sigma^{-1}\}$ is bounded, we are led to take $\sigma \sim R$; in this case the term $e^{\frac{\sigma^2}{2} \sup_j \Im m(r_j)^2}$ will be too large to yield any interesting information.

We see that the method only has a chance to work if $T\Theta(R) \gg \log R$. From now on we take $\Theta(R) \geq \log R$, and always such that $R^{-1}\Theta(R) \rightarrow 0$. We also take $\sigma^{-2} = C\Theta(R)$ with C large. We must note that the parameters r_j correspond to the eigenvalues of $-\Delta_{\Theta(R)\omega}$, and thus they also depend on R .

8.3 Proof of Theorem 1.12. The right-hand side of (8.3) is easy to understand. The term

$$\int_{-\infty}^{+\infty} r \tanh \pi r \left[e^{-\frac{\sigma^2}{2}(r-R)^2} e^{-iTr} + e^{-\frac{\sigma^2}{2}(r+R)^2} e^{iTr} \right] dr \tag{8.16}$$

is $\mathcal{O}(\sigma^{-1}R)$, whereas the \sum_{γ} has modulus greater than

$$\begin{aligned}
 & \frac{1}{2} \sum_{\ell_\gamma \leq 5T} \frac{e^{\int_\gamma \omega} \ell_{\gamma_0}}{\sinh \frac{\ell_\gamma}{2}} \frac{1}{\sqrt{2\pi\sigma}} \left[e^{-\frac{(\ell_\gamma - T)^2}{2\sigma^2}} + e^{-\frac{(\ell_\gamma + T)^2}{2\sigma^2}} \right] \\
 & \quad + \sum_{\ell_\gamma \geq 5T} \frac{e^{\int_\gamma \omega} \ell_{\gamma_0}}{\sinh \frac{\ell_\gamma}{2}} \frac{1}{\sqrt{2\pi\sigma}} \left[e^{-\frac{(\ell_\gamma - T)^2}{2\sigma^2}} + e^{-\frac{(\ell_\gamma + T)^2}{2\sigma^2}} \right] \cos(\ell_\gamma R) \\
 & \geq \frac{1}{2} \sum_{T-1 \leq \ell_\gamma \leq T} \frac{e^{\int_\gamma \omega} \ell_{\gamma_0}}{\sinh \frac{\ell_\gamma}{2}} \frac{1}{\sqrt{2\pi\sigma}} \left[e^{-\frac{(\ell_\gamma - T)^2}{2\sigma^2}} + e^{-\frac{(\ell_\gamma + T)^2}{2\sigma^2}} \right] \\
 & \quad + \sum_{\ell_\gamma \geq 5T} \frac{e^{\int_\gamma \omega} \ell_{\gamma_0}}{\sinh \frac{\ell_\gamma}{2}} \frac{1}{\sqrt{2\pi\sigma}} \left[e^{-\frac{(\ell_\gamma - T)^2}{2\sigma^2}} + e^{-\frac{(\ell_\gamma + T)^2}{2\sigma^2}} \right] \cos(\ell_\gamma R) \\
 & \geq \frac{1}{\sqrt{8\pi\sigma}} \left[e^{TH(1-\delta) - T/2} e^{\Theta(R)(1-\delta)T} e^{-1/2\sigma^2} + \mathcal{O}(1) \right], \tag{8.17}
 \end{aligned}$$

(using (8.12)), and thus is much greater than the integral (8.16). To get the last $\mathcal{O}(1)$ we have used the exponential growth of the number of closed geodesics. The left-hand side of (8.3) is more complicated to bound from above, since the r_j now depend on R .

PROPOSITION 5. *Take $\Theta(R) \geq \log R$, and such that $R^{-1}\Theta(R) \rightarrow 0$. Take $\sigma^{-2} = C\Theta(R)$. Let $f(R)$ be such that $\sigma^2 f(R)^2 \gg T\Theta(R)$. Then*

$$\begin{aligned}
 \left| \sum_j e^{-\frac{\sigma^2}{2}(r_j - R)^2} e^{-iTr_j} \right| & \leq \#\{j, |\Re e(r_j) - R| \\
 & \leq f(R)\} e^{\frac{\sigma^2}{2} \sup_j \Im m(r_j)^2} e^{T \sup_j |\Im m(r_j)|} + \mathcal{O}(1),
 \end{aligned}$$

where the \sup_j are taken over the set of indices j such that $|\Re e(r_j) - R| \leq f(R)$.

Proof.

LEMMA 8.3. *We have an a priori bound $|\Im m(r_j)| \leq c\Theta(R)$, where c depends only on ω_o .*

Indeed, let $r_j = x + iy$ and $f \in L^2(M)$ be such that $\|f\|_{L^2} = 1$ and

$$-\Delta f + 2\Theta(R)\langle \omega_o, df \rangle - \Theta(R)^2 \|\omega_o\|_x^2 f = \left(\frac{1}{4} + r_j^2\right) f.$$

Taking the scalar product with f , we get

$$\frac{1}{4} + x^2 - y^2 = \int |\nabla f|^2 - \Theta(R)^2 \int \|\omega_o\|_x^2 |f|^2 \tag{8.18}$$

and

$$2xy = -2\Theta(R)i \int \langle \omega_o, df \rangle \bar{f}. \tag{8.19}$$

Equation (8.19) yields $|xy| \leq c\Theta(R)\sqrt{\int |\nabla f|^2}$. Equation (8.18) implies that $x^2 \geq \int |\nabla f|^2 - c^2\Theta(R)^2$. If $|x| \geq \frac{1}{2}\sqrt{\int |\nabla f|^2}$ then we are done, by (8.19). If $|x| \leq \frac{1}{2}\sqrt{\int |\nabla f|^2}$, then (8.18) implies that $\int |\nabla f|^2 \leq 2c^2\Theta(R)^2$ and that $y^2 \leq \frac{1}{4} + 5c^2\Theta(R)^2$. The lemma follows.

We now break the sum $\sum_j e^{-\frac{\sigma^2}{2}(r_j-R)^2} e^{-iT r_j}$ into three parts: $I = \sum_{j, \Re(r_j) \leq R-f(R)}$, $II = \sum_{j, |\Re(r_j) - R| \leq f(R)}$ and $III = \sum_{j, \Re(r_j) \geq R+f(R)}$.

The last sum III is bounded by

$$e^{\frac{\sigma^2}{2} c^2 \Theta(R)^2} e^{cT\Theta(R)} \sum_{j, \Re(r_j) \geq R+f(R)} e^{-\frac{\sigma^2}{2} (\Re(r_j) - R)^2}.$$

We decompose this sum into $\sum_{n \geq 0} \sum_{R+f(R)+n \leq \Re(r_j) \leq R+f(R)+n+1}$, and by Weyl's law in the form (8.8), this is dominated by

$$e^{cT\Theta(R)} \sum_{n \geq 0} (R+n+f(R)) \Theta(R) e^{-\frac{\sigma^2}{2} (n+f(R))^2} \leq e^{cT\Theta(R)} \Theta(R) \int_{f(R)-1}^{+\infty} (R+x) e^{-\frac{\sigma^2 x^2}{2}} dx$$

and with our relations between $T, R, f(R)$ and σ , this last quantity is $\mathcal{O}(1)$.

Concerning the first sum I , we bound it by

$$e^{\frac{\sigma^2}{2} c^2 \Theta(R)^2} e^{cT\Theta(R)} \sum_{j, \Re(r_j) \leq R-f(R)} e^{-\frac{\sigma^2}{2} (\Re(r_j) - R)^2}. \tag{8.20}$$

The subsum $\sum_{j, \Re(r_j) \leq -R+f(R)}$ can be treated as above and shown to be $\mathcal{O}(1)$ (using Weyl's law in the form (8.8)), and we only need to concentrate on $\sum_{j, |\Re(r_j)| \leq R-f(R)}$. To bound this sum, we first need a control the number of terms.

LEMMA 8.4. $\#\{j, |\Re(r_j)| \leq R\} = \mathcal{O}(R^2 \Theta(R))$.

To that end, we use again the trace formula and write

$$\sum_j e^{-r_j^2/2R^2} = \frac{\text{Area}(M)}{4\pi} \int_{-\infty}^{+\infty} r \tanh(\pi r) \cdot e^{-r^2/2R^2} dr + \sum_{\gamma} \frac{e^{\Theta(R) \int_{\gamma} \omega_{\alpha}} \ell_{\gamma_{\alpha}}}{\sinh \ell_{\gamma}/2} \frac{R}{\sqrt{2\pi}} e^{-R^2 \ell_{\gamma}^2/2}.$$

On the right, the term \sum_{γ} is clearly $o(1)$ whereas the \int term is of order R^2 . On the left, we break the sum into $\sum_{j, |\Re(r_j)| \leq R}$ and $\sum_{j, |\Re(r_j)| \geq R}$. As above, we can use Weyl's law in the form (8.8) to show that the $\sum_{j, |\Re(r_j)| \geq R} e^{-r_j^2/2R^2}$ is $\mathcal{O} \sum_{n \geq 0} (R+n+1) \Theta(R) e^{-(R+n)^2/2R^2} = \mathcal{O}(R^2 \Theta(R))$. Thus we have

$$\sum_{j, |\Re(r_j)| \leq R} e^{-\frac{\Re(r_j)^2 - \Im m(r_j)^2}{2R^2}} e^{-\frac{i \Re(r_j) \Im m(r_j)}{R^2}} = \mathcal{O}(R^2 \Theta(R)).$$

Remember that $|\Im m(r_j)| \leq c \Theta(R)$ and that $R^{-1} \Theta(R) \rightarrow 0$. Thus, for $|\Re(r_j)| \leq R$ we can write $e^{-i \Re(r_j) \Im m(r_j)/R^2} = 1 + \mathcal{O}(R^{-1} \Theta(R))$. This yields

$$\begin{aligned} e^{-1/2} \#\{j, |\Re(r_j)| \leq R\} (1 + \mathcal{O}(R^{-1} \Theta(R))) & \\ & \leq \sum_{j, |\Re(r_j)| \leq R} e^{-\frac{\Re(r_j)^2 - \Im m(r_j)^2}{2R^2}} (1 + \mathcal{O}(R^{-1} \Theta(R))) \\ & = \mathcal{O}(R^2 \Theta(R)), \end{aligned}$$

and finishes the proof of Lemma 8.4.

Now, we go back to the sum $\sum_{j, |\Re(r_j)| \leq R-f(R)}$ in (8.20), and we see that it is bounded by

$$R^2 \Theta(R) e^{\frac{\sigma^2}{2} c^2 \Theta(R)^2} e^{cT\Theta(R)} e^{-\sigma^2 f(R)^2/2} = \mathcal{O}(1).$$

This ends the proof of Proposition 5 □

We can now come back to (8.3). Noting that $\sigma^2 \sup_j |\Im m(r_j)|^2 = \mathcal{O}(1)$, Proposition 5 shows that the left-hand side of the trace formula (8.3) is bounded from above by

$$C \#\{j, |\Re e(r_j) - R| \leq f(R)\} e^{T \sup_j |\Im m(r_j)|} + \mathcal{O}(1) \leq CRf(R)\Theta(R)e^{T \sup_j |\Im m(r_j)|},$$

where the sup is taken over all j such that $|\Re e(r_j) - R| \leq f(R)$, and where we have used again (8.8). We can of course assume, without loss of generality, that $f(R) = \mathcal{O}(R)$.

On the right-hand side of (8.3), the f is $\mathcal{O}(R\Theta(R)^{1/2})$, and the \sum_γ is bounded from below by

$$\frac{1}{\sqrt{8\pi\sigma}} e^{TH(1-\delta)-T/2} e^{\Theta(R)(1-\delta)T} e^{-\frac{1}{2\sigma^2}},$$

as was proved in (8.17). Writing

$$CRf(R)\Theta(R)e^{T \sup_j |\Im m(r_j)|} \geq \sigma^{-1} e^{TH(1-\delta)-T/2} e^{\Theta(R)(1-\delta)T} e^{-1/2\sigma^2},$$

remembering that $\Theta(R) \geq \log(R)$, $\Theta(R) = o(R)$, $\sigma^{-2} = c\Theta(R)$ and $T \asymp \log \log R$, we see that necessarily

$$\sup_j |\Im m(r_j)| \geq (1 - 2\delta)\Theta(R),$$

where the sup is taken over all j such that $|\Re e(r_j) - R| \leq f(R)$. This finishes the proof of Theorem 1.12.

9 The Arithmetic Case

We now prove Theorem 1.9, and first we need to define the special class of hyperbolic surfaces we are considering. Let $p \geq 3$ be a prime, $p \equiv 1 \pmod{4}$, and $A \geq 1$ be a quadratic non-residue modulo p . We set

$$\Gamma = \Gamma(A, p) = \left\{ \begin{pmatrix} y_0 + y_1 \sqrt{A} & y_2 \sqrt{p} + y_3 \sqrt{Ap} \\ y_2 \sqrt{p} - y_3 \sqrt{Ap} & y_0 - y_1 \sqrt{A} \end{pmatrix} \in SL(2, \mathbb{R}), y_0, y_1, y_2, y_3 \in \mathbb{Z} \right\}.$$

It is a discrete cocompact subgroup of $SL(2, \mathbb{R})$ which contains only hyperbolic transformations [H]. We consider the hyperbolic surface $M = \Gamma \backslash \mathbb{H}$. In M , the lengths of the closed geodesics are the $\log x_m$, where

$$x_m = 2m^2 - 1 + 2m\sqrt{m^2 - 1}, \quad m \in \mathbb{N}.$$

We define

$$\mu(m) = \sum_{\gamma, \ell_\gamma = \log x_m} e^{J_\gamma \omega} \ell_{\gamma_0}.$$

We now follow very closely the approach of [H, p.304-314]. We introduce an even function k on \mathbb{R} , whose Fourier transform is non-negative and compactly supported in $[-1, 1]$; we also assume that $\hat{k} \geq 1$ on $[-1/2, 1/2]$. We define

$$K_\alpha(r) = k(r)[e^{i\alpha r} + e^{-i\alpha r}].$$

We write again the trace formula: for all $t > 0$,

$$\sum_j K_\alpha(r_j - t) + K_\alpha(r_j + t) = \frac{\text{Area}(M)}{4\pi} \int r \tanh(\pi r) [K_\alpha(r - t) + K_\alpha(r + t)] dr + 2 \sum_\gamma \frac{e^{\int_\gamma \omega} \ell_{\gamma_0}}{\sinh \ell_\gamma / 2} \hat{K}_\alpha(\ell_\gamma) \cos(t \ell_\gamma). \quad (9.1)$$

We will bound from below the right-hand side (averaged in t) to obtain information on the left-hand side. Denote

$$\begin{aligned} S_\alpha(t) &= \sum_\gamma \frac{e^{\int_\gamma \omega} \ell_{\gamma_0}}{\sinh \ell_\gamma / 2} \hat{K}_\alpha(\ell_\gamma) \cos(t \ell_\gamma) \\ &= 2 \sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \frac{\mu(m)}{x_m^{1/2} - x_m^{-1/2}} \hat{K}_\alpha(\log x_m) \cos(t \log x_m). \end{aligned}$$

PROPOSITION 6. *Let $\alpha = 2\beta \log T - C$, $0 < \beta \leq 1$ and C large enough. Then,*

$$\int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) |S_\alpha(t)|^2 dt \geq \tilde{C} T^{\beta(4\text{Pr}(\omega)-2)}.$$

Although we are actually interested in the quantity $\int_{2T-T^\beta}^{2T+T^\beta} |S_\alpha(t)|^2 dt$, the reason for introducing the regularizing factor $(1 - \frac{|t-2T|}{T^\beta})$ is exactly the same as in [H, p. 315], we want a sufficiently regular function, for Lemma 9.1 to hold.

Proof. Introduce the notation

$$\begin{aligned} \eta(m) &= \frac{\mu(m)}{x_m^{1/2} - x_m^{-1/2}} \hat{K}_\alpha(\log x_m), \\ \nu(m) &= \mu(m) \hat{K}_\alpha(\log x_m). \end{aligned}$$

Divide the integral $I = \int_{2T-T^\beta}^{2T+T^\beta} (1 - \frac{|t-2T|}{T^\beta}) |S_\alpha(t)|^2 dt$ into $I = I_1 + I_2$, where

$$I_1 = \sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \eta(m)^2 \int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) \cos^2(t \log x_m) dt$$

and

$$I_2 = 2 \sum_{e^{\alpha-1} \leq x_k < x_m \leq e^{\alpha+1}} \eta(m)\eta(k) \int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) \cos(t \log x_m) \cos(t \log x_k) dt.$$

The idea is that the x_m, x_k , with $x_m \neq x_k$, are well-spaced, implying that the oscillatory integral I_2 is small compared to I_1 .

LEMMA 9.1. *For $T \geq 1$ and $\lambda \in \mathbb{R}$,*

$$\int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) e^{i\lambda t} dt = \mathcal{O} \left[\min \left(T^\beta, \frac{1}{\lambda^2 T^\beta} \right) \right].$$

Let us first consider I_1 .

$$\int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) \cos^2(t \log x_m) dt = \int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) \frac{1 + \cos(2t \log x_m)}{2} dt$$

$$= \frac{T^\beta}{2} + \mathcal{O}(T^{-\beta} \log x_m^{-2}).$$

Hence,

$$I_1 \geq \left(\frac{T^\beta}{2} + o(1) \right) \sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \eta(m)^2 \tag{9.2}$$

$$\geq \left(\frac{T^\beta}{2} + o(1) \right) \sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \frac{\nu(m)^2}{(x_m^{1/2} - x_m^{-1/2})^2} \tag{9.3}$$

$$\geq c_1 T^\beta e^{-\alpha} \sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \nu(m)^2. \tag{9.4}$$

The right-hand side of (9.4) will be estimated later. We now turn to I_2 , and want to show that it is much smaller than I_1 . The integral

$$\int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta} \right) \cos(t \log x_m) \cos(t \log x_k) dt$$

is smaller than $\frac{1}{T^\beta (\log x_m - \log x_k)^2}$. We can ensure that

$$\frac{1}{|\log x_m - \log x_k|} \leq cT^\beta$$

for $e^{\alpha-1} < x_k < x_m < e^{\alpha+1}$, by choosing α in an appropriate range: writing

$$\log x_m - \log x_{m-1} \sim \frac{x_m - x_{m-1}}{x_m} \sim \frac{2}{m} \sim \frac{4}{\sqrt{x_m}},$$

we see we have to take $\alpha \leq 2\beta \log T - C$ (C large). More generally, we have by the intermediate value theorem

$$|\log x_m - \log x_k| \geq \tilde{C} e^{-\alpha/2} |m - k|.$$

The analysis done by [H, p. 310-311] can be applied *verbatim* to show that

$$|I_2| \leq \frac{\tilde{C}}{T^\beta} \sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \nu(m)^2. \tag{9.5}$$

Comparing (9.4) and (9.5), we see that

$$|I_2| \leq \frac{I_1}{100},$$

provided $\alpha \leq 2\beta \log T - C$ with C sufficiently large. Thus,

$$I \geq \frac{99}{100} I_1.$$

To complete our estimate for I , we must return to equation (9.4). Clearly,

$$\sum_{e^{\alpha-1} \leq x_m \leq e^{\alpha+1}} \nu(m)^2 \geq \sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} \mu(m)^2.$$

We write the Cauchy–Schwarz inequality,

$$\left[\sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} \mu(m) \right]^2 \leq \left(\sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} 1 \right) \left(\sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} \mu(m)^2 \right).$$

But

$$\sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} \mu(m) = \sum_{\gamma, \alpha-1/2 \leq \ell_\gamma \leq \alpha+1/2} e^{\int_\gamma \omega} \ell_{\gamma_0} \geq \tilde{C} e^{\alpha \text{Pr}(\omega)},$$

see [PP, p. 117]. On the other hand,

$$\sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} 1 = \mathcal{O}(e^{\alpha/2}).$$

We obtain this way

$$\sum_{e^{\alpha-1/2} \leq x_m \leq e^{\alpha+1/2}} \mu(m)^2 \geq C e^{2\alpha \text{Pr}(\omega) - \alpha/2}.$$

We have proved

$$\int_{2T-T^\beta}^{2T+T^\beta} \left(1 - \frac{|t-2T|}{T^\beta}\right) |S_\alpha(t)|^2 dt \geq \tilde{C} T^{\beta(4\text{Pr}(\omega)-2)}. \quad \square$$

This implies

$$|S_\alpha(t)| \geq C t^{\beta(2\text{Pr}(\omega) - \frac{3}{2})}$$

for some $t \in [2T - T^\beta, 2T + T^\beta]$.

Now consider the integral $\int r \tanh(\pi r) K_\alpha(r-t) dr$ or $\int r \tanh(\pi r) K_\alpha(r+t) dr$ in (9.1). We write

$$\int r \tanh(\pi r) K_\alpha(r-t) dr = \int r \tanh(\pi r) k(r-t) [e^{i\alpha(r-t)} + e^{-i\alpha(r-t)}] dr.$$

To evaluate $\int (r+t) \tanh(\pi(r+t)) k(r) e^{i\alpha r} dr$, we shift the integral over \mathbb{R} to an integral over $(\frac{1}{2} - \epsilon)i + \mathbb{R}$, and we find that the integral is $\mathcal{O}(te^{-|\alpha|(\frac{1}{2}-\epsilon)})$ for any $\epsilon > 0$.

LEMMA 9.2. $\int r \tanh(\pi r) K_\alpha(r-t) dr$ or $\int r \tanh(\pi r) K_\alpha(r+t) dr = \mathcal{O}(te^{-|\alpha|(\frac{1}{2}-\epsilon)}) = \mathcal{O}(t^{1+\epsilon-\beta})$ for any $\epsilon > 0$.

We finally turn to $\sum_j K_\alpha(r_j - t) + K_\alpha(r_j + t)$. Fixing a small $\epsilon > 0$, one sees using Weyl's law, the fact that k is rapidly decreasing in any horizontal strip – and the fact that the $\Im m(r_j)$ are bounded (by $c\|\omega\|$) – that

$$\sum_j K_\alpha(r_j - t) = \sum_{j, |\Re(r_j) - t| \leq t^\epsilon} K_\alpha(r_j - t) + \mathcal{O}(t^{-\infty}).$$

Similarly,

$$\sum_j K_\alpha(r_j + t) = \sum_{j, |\Re(r_j) + t| \leq t^\epsilon} K_\alpha(r_j + t) + \mathcal{O}(t^{-\infty}).$$

We see that

$$\left| \sum_j K_\alpha(r_j - t) + K_\alpha(r_j + t) \right| \leq C e^{\alpha \sup_j |\Im m(r_j)| t^{1+\epsilon}} \leq C t^{2\beta \sup_j |\Im m(r_j)| t^{1+\epsilon}}$$

where the \sup_j is taken over the j such that $|\Re(r_j) \pm t| \leq t^\epsilon$. Here the factor C is essentially $\sup\{|k(z)|, |\Im m(z)| \leq c\|\omega\|\}$, it does not depend on t .

We have proved that there exists $t \in [T - T^\beta, T + T^\beta]$, and r_j with $|\Re(r_j) - t| \leq t^\epsilon$, such that

$$t^{2\beta \sup_j |\Im(r_j)|} t^{1+\epsilon} \geq \tilde{C} t^{\beta(2\Pr(\omega) - \frac{3}{2})}.$$

In particular, if T is large enough, this implies

$$\sup \{ |\Im(r_j)|, |\Re(r_j) \pm T| \leq T^\beta \} \geq \Pr(\omega) - \frac{3}{4} - \frac{1 + \epsilon}{2\beta},$$

and this proves Theorem 1.9.

10 Appendix: Symbol Classes

Following [DS], for any $0 \leq \delta < 1/2$ we introduce the symbol class

$$S_\delta^m \stackrel{\text{def}}{=} \{ a \in C^\infty(T^*M), |\partial_x^\alpha \partial_\xi^\beta a| \leq C_{\alpha,\beta} \hbar^{-\delta|\alpha+\beta|} \langle \xi \rangle^{m-|\beta|} \}. \quad (10.1)$$

If $\delta = 0$ we will just denote S^m . We will denote $\text{Op}_\hbar(a(x, \xi)) = \text{Op}(a(x, \hbar\xi))$, where Op is a quantization procedure on M . The quantization of any $a \in S_\delta^0$ leads to a bounded operator on $L^2(M)$ (the norm being bounded uniformly in \hbar), see [DS]. We will denote $\Psi DO_\delta^m = \text{Op}_\hbar(S_\delta^m)$.

We use

PROPOSITION 7. *Let $a \in S_\delta^m, b \in S_{\delta'}^n$ with $0 \leq \delta' \leq \delta < 1/2$. Then,*

- (i) $\text{Op}_\hbar(a) \text{Op}_\hbar(b) - \text{Op}_\hbar(ab) \in \hbar^{1-\delta-\delta'} \text{Op}_\hbar(S_\delta^{m+n-1})$.
- (ii) $[\text{Op}_\hbar(a), \text{Op}_\hbar(b)] - \frac{\hbar}{i} \text{Op}_\hbar(\{a, b\}) \in \hbar^{2(1-\delta-\delta')} \text{Op}_\hbar(S_\delta^{m+n-2})$.

We also use a local form of the Calderon–Vaillancourt estimate [DS]:

PROPOSITION 8. *There exists $K \in \mathbb{N}$ depending only on the dimension of M , such that the following holds. Take $A = \text{Op}_\hbar(a)$ where $a \in S_\delta^2$. Let I be an open interval of \mathbb{R}^+ and let λ belong to I . Then, there exists $C > 0$, and $C(a, \lambda)$ depending on a finite number of seminorms of a (uniform in λ if it stays inside a compact subset of I), such that, for all $u \in L^2(M)$,*

$$\|Au\|_{L^2} \leq C \left(\sup_{p_o^{-1}(I)} |a| + \sum_{k=1}^K \hbar^k \sup_{p_o^{-1}(I)} |D^{2k}a| \right) \|u\|_{L^2} + C(a, \lambda) \|(P - \lambda)u\|_{L^2}.$$

In fact $C(a, \lambda)$ is controlled by the supremum norm of $\frac{a}{p_o^{-\lambda}}$ and a finite number of its derivatives outside $p_o^{-1}(I)$. Similarly we have

$$|\langle u, Au \rangle| \leq C \left(\sup_{p_o^{-1}(I)} |a| + \sum_{k=1}^K \hbar^k \sup_{p_o^{-1}(I)} |D^{2k}a| \right) \|u\|_{L^2} + C(a, \lambda) \|(P - \lambda)u\|_{L^2}^2.$$

11 Appendix: Sjöstrand’s Proof of Proposition 3

In order to prove Proposition 3, we first want to bound the norm and trace norm of

$$f \left(\frac{2P - 1}{\hbar} \right) (\tilde{Q}_T - Q_T) f \left(\frac{2P - 1}{\hbar} \right).$$

We write a Calderon–Vaillancourt type estimate,

$$\|(\tilde{Q}_T - Q_T)u\| \leq \left(\sup_{p_o^{-1}[\frac{1}{2}-\epsilon, \frac{1}{2}+\epsilon]} (\tilde{q}^T - q^T) + \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\| + \mathcal{O}(1) \|(2P - 1)u\|,$$

where δ is as in Proposition 2. Besides, $\|(2P - 1)f\left(\frac{2P-1}{\hbar}\right)\| = \mathcal{O}(\hbar)$. It follows that

$$\left\| f\left(\frac{2P-1}{\hbar}\right) (\tilde{Q}_T - Q_T) f\left(\frac{2P-1}{\hbar}\right) \right\| \leq \|f\|_\infty^2 \left(\sup_{p_o^{-1}[\frac{1}{2}-\epsilon, \frac{1}{2}+\epsilon]} (\tilde{q}^T - q^T) + \mathcal{O}(\hbar^{1-2\delta}) \right).$$

For the trace class norm, we need to be even more careful than in [Sj]. Instead of using the Gårding inequality, we use the existence of a positive quantization – meaning that $\text{Op}_\hbar(a) \geq 0$ if $a \geq 0$. If we choose such, we have directly that $f\left(\frac{2P-1}{\hbar}\right) (\tilde{Q}_T - Q_T) f\left(\frac{2P-1}{\hbar}\right) \geq 0$ in the operator sense. Thus,

$$\begin{aligned} \left\| f\left(\frac{2P-1}{\hbar}\right) (\tilde{Q}_T - Q_T) f\left(\frac{2P-1}{\hbar}\right) \right\|_1 &= \text{Tr} f\left(\frac{2P-1}{\hbar}\right) (\tilde{Q}_T - Q_T) f\left(\frac{2P-1}{\hbar}\right) \\ &= \text{Tr} f\left(\frac{2P-1}{\hbar}\right)^2 (\tilde{Q}_T - Q_T) \\ &= \text{Tr} \frac{1}{2\pi} \int \hat{f}^2(t) e^{it\frac{2P-1}{\hbar}} (\tilde{Q}_T - Q_T) dt. \end{aligned}$$

Writing the expansion of $e^{it\frac{2P-1}{\hbar}}$ as a Fourier integral operator, writing the trace as the integral of the kernel, and applying the stationary phase method in the time-energy variables, we obtain an asymptotic expansion

$$\begin{aligned} \text{Tr} \frac{1}{2\pi} \int \hat{f}^2(t) e^{it\frac{2P-1}{\hbar}} (\tilde{Q}_T - Q_T) dt \\ = C_d \hbar^{2-d} \left[\hat{f}^2(0) \int_{p_o^{-1}(1/2)} (q^T - \tilde{q}^T) L_{1/2}(d\rho) \right. \\ \left. + \sum_{k=1}^{N-1} \hbar^k D_t^{2k} \hat{f}^2(0) \int_{p_o^{-1}(1/2)} D_\rho^{2k} (q^T - \tilde{q}^T) L_{1/2}(d\rho) + \mathcal{O}(\hbar^{N(1-2\delta)}) \right], \end{aligned}$$

where D_t^{2k} and D_ρ^{2k} are differential operators of degree $\leq 2k$, respectively on \mathbb{R} and T^*M . Note that the term $\hbar^k D_t^{2k} \hat{f}^2(0) \int_{p_o^{-1}(1/2)} D_\rho^{2k} (q^T - \tilde{q}^T) L_{1/2}(d\rho)$ is a $\mathcal{O}(\hbar^{k(1-2\delta)}) L_{1/2}(\tilde{q}^T \neq q^T) = o(1) L_{1/2}(\tilde{q}^T \neq q^T)$. This proves, in particular, Corollary 3.

To finish the proof of Proposition 3, there remains to study the invertibility of $z - \tilde{\mathcal{P}}_T$. Recall the identity

$$\|(A + iB)u\|^2 = \|Au\|^2 + \|Bu\|^2 + i\langle u, [A, B]u \rangle,$$

if A, B are bounded self-adjoint operators. Thus,

$$\begin{aligned} 2\|(\tilde{\mathcal{P}}_T - z)u\|^2 &\geq \|(P + i\hbar\hat{Q}_T - z)u\|^2 - \mathcal{O}(\hbar^{4(1-\delta)}) (\|(2P - 1)u\|^2 + \|u\|^2) \\ &\geq \|(P - \Re e(z))u\|^2 + \hbar^2 \left\| \left(\frac{\Im m(z)}{\hbar} - \hat{Q}_T \right) u \right\|^2 + i\hbar \langle u, [P, \hat{Q}_T]u \rangle \\ &\quad - \mathcal{O}(\hbar^{4(1-\delta)}) (\|(2P - 1)u\|^2 + \|u\|^2) \end{aligned}$$

$$\begin{aligned}
 &= \|(P - \Re e(z))u\|^2 + \hbar^2 \left\| \left(\frac{\Im m(z)}{\hbar} - \hat{Q}_T \right) u \right\|^2 \\
 &\quad + \left(\mathcal{O}(1) \frac{\hbar^2}{T} (1 + \|f\|_\infty^2) + \mathcal{O}(\hbar^{3-2\delta}) \right) \|u\|^2 + \mathcal{O}(\hbar^2) \|(P - \Re e(z))u\|^2.
 \end{aligned}$$

We have used (4.3) (or Proposition 8), and the same for \tilde{Q}_T . We find that

$$\sqrt{3} \|(\tilde{P}_T - z)u\| \geq \|(P - \Re e(z))u\| - \left(\mathcal{O} \left(\frac{\hbar}{\sqrt{T}} \right) + \mathcal{O}(\hbar^{\frac{3}{2}-\delta}) \right) \|u\|. \tag{11.1}$$

On the other hand, we have

$$\begin{aligned}
 &\Im m \left\langle \frac{1}{\hbar} (z - \tilde{P}_T)u, u \right\rangle \\
 &= \left\langle \left(\frac{\Im m(z)}{\hbar} - \hat{Q}_T \right) u, u \right\rangle + \mathcal{O}(\hbar^{1-2\delta}) (\|u\| + \|(\Re e(z) - P)u\|) \|u\| \\
 &\quad \left\langle \left(\frac{\Im m(z)}{\hbar} - Q_T + f \left(\frac{2P - 1}{\hbar} \right) (\tilde{Q}_T - Q_T) f \left(\frac{2P - 1}{\hbar} \right) \right) u, u \right\rangle \\
 &\quad + \mathcal{O}(\hbar^{1-2\delta}) (\|u\| + \|(\Re e(z) - P)u\|) \|u\| \\
 &= \left\langle \left(\frac{\Im m(z)}{\hbar} - Q_T + f \left(\frac{2\Re e(z) - 1}{\hbar} \right)^2 (\tilde{Q}_T - Q_T) \right) u, u \right\rangle \\
 &\quad + \mathcal{O}(\hbar^{1-2\delta}) (\|u\| + \|(\Re e(z) - P)u\|) \|u\| \\
 &\quad - \left(2 \sup_{p_o^{-1}[\frac{1}{2}-\epsilon, \frac{1}{2}+\epsilon]} (q^T - \tilde{q}^T) \|f\|_\infty \|f'\|_\infty + \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\| \left\| \frac{P - \Re e(z)}{\hbar} u \right\|
 \end{aligned} \tag{11.2}$$

by the same trick as in [Sj, (3.19)]. Recall that we are interested in a region where $z - \frac{1}{2} = \mathcal{O}(\hbar)$.

Let $F(E) > 0$ be a continuous function defined on a bounded interval J containing 0, and restrict z by assuming that

$$\frac{\Im m(\zeta)}{2\hbar} - q^T + f \left(\frac{\Re e(\zeta)}{\hbar} \right)^2 (q^T - \tilde{q}^T) \geq F \left(\frac{\Re e(\zeta)}{\hbar} \right),$$

near $p_o^{-1}(1/2)$, $\Re e(\zeta)/\hbar \in J$ (where $\zeta = 2z - 1$). It follows from the Gårding inequality that for such z

$$\begin{aligned}
 &\left\langle \left(\frac{\Im m(z)}{\hbar} - Q_T + f \left(\frac{2\Re e(z) - 1}{\hbar} \right)^2 (\tilde{Q}_T - Q_T) \right) u, u \right\rangle \\
 &\quad \geq \left(F \left(\frac{\Re e(\zeta)}{\hbar} \right) - \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\|^2 - \mathcal{O}(1) \|u\| \|(P - \Re e(z))u\|.
 \end{aligned}$$

Using this in (11.2), we get

$$\begin{aligned}
 &\Im m \left\langle \frac{1}{\hbar} (z - \tilde{P}_T)u, u \right\rangle \geq \left(F \left(\frac{\Re e(\zeta)}{\hbar} \right) - \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\|^2 \\
 &\quad - \left(2 \sup_{p_o^{-1}[\frac{1}{2}-\epsilon, \frac{1}{2}+\epsilon]} (q^T - \tilde{q}^T) \|f\|_\infty \|f'\|_\infty + \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\| \left\| \frac{P - \Re e(z)}{\hbar} u \right\|. \tag{11.3}
 \end{aligned}$$

Reasoning as in [Sj], (3.25) and (3.26), we find finally

$$\begin{aligned} \Im m \left\langle \frac{1}{\hbar} (z - \tilde{\mathcal{P}}_T) u, u \right\rangle &\geq \left(F \left(\frac{\Re e(\zeta)}{\hbar} \right) - \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\|^2 \\ &\quad - \sqrt{3} \left(2 \sup_{p\sigma^{-1}]_{\frac{1}{2}-\epsilon, \frac{1}{2}+\epsilon[}} (q^T - \tilde{q}^T) \|f\|_\infty \|f'\|_\infty + \mathcal{O}(\hbar^{1-2\delta}) \right) \|u\| \left\| \frac{\tilde{\mathcal{P}}_T - z}{\hbar} u \right\| \\ &\quad - \left(\mathcal{O} \left(\frac{1}{\sqrt{T}} \right) + \mathcal{O}(\hbar^{\frac{1}{2}-\delta}) \right) \|u\|^2. \end{aligned}$$

and

$$\begin{aligned} &\left(F \left(\frac{\Re e(\zeta)}{\hbar} \right) + \mathcal{O} \left(\frac{1}{\sqrt{T}} \right) + \mathcal{O}(\hbar^{\frac{1}{2}-\delta}) \right) \|u\| \\ &\leq \left[1 + 2\sqrt{3} \sup_{p\sigma^{-1}]_{\frac{1}{2}-\epsilon, \frac{1}{2}+\epsilon[}} (q^T - \tilde{q}^T) \|f\|_\infty \|f'\|_\infty + \mathcal{O}(\hbar^{1-2\delta}) \right] \left\| \frac{\tilde{\mathcal{P}}_T - z}{\hbar} u \right\| \end{aligned}$$

which finishes the proof of Proposition 3.

References

- [A1] N. ANANTHARAMAN, Precise counting results for closed orbits of Anosov flows, *Ann. Sci. École Norm. Sup.* (4) 33:1 (2000), 33–56.
- [A2] N. ANANTHARAMAN, Counting geodesics which are optimal in homology, *Ergodic Theory Dynam. Systems* 23:2 (2003), 353–388.
- [An] D.V. ANOSOV, Geodesic Flows on Closed Riemannian Manifolds of Negative Curvature, *Trudy Mat. Inst. Steklov.* 90, 1967.
- [BL] M. BABILLOT, M. LEDRAPPIER, Lalley’s theorem on periodic orbits of hyperbolic flows, *Ergodic Theory Dynam. Systems* 18:1 (1998), 17–39.
- [BR] R. BOWEN, D RUELLE, The ergodic theory of Axiom A flows, *Invent. Math.* 29:3 (1975), 181–202.
- [DS] M. DIMASSI, J. SJÖSTRAND, Spectral Asymptotics in the Semi-classical Limit, London Mathematical Society Lecture Note Series 268, Cambridge University Press, Cambridge (1999).
- [GK] V. GUILLEMIN, D. KAZHDAN, On the cohomology of certain dynamical systems, *Topology* 19:3 (1980), 291–299.
- [H] D.A. HEJHAL, The Selberg Trace Formula for $\mathrm{PSL}(2, R)$, Vol. I, Springer Lecture Notes in Mathematics 548 (1976).
- [HiSV] M. HITRIK, J. SJÖSTRAND, S. VÛ NGỌC, Diophantine tori and spectral asymptotics for nonselfadjoint operators, *Amer. J. Math.* 129:1 (2007), 105–182.
- [JN] D. JAKOBSON, F. NAUD, Lower bounds for resonances of infinite area Riemann surfaces, *Analysis & PDE* 3:2 (2010), to appear.
- [JP] D. JAKOBSON, I. POLTEROVICH, Estimates from below for the spectral function and for the remainder in local Weyl’s law, *Geom. Funct. Anal.* 17:3 (2007), 806–838.
- [KH] A. KATOK, B. HASSELBLATT, Introduction to the modern theory of dynamical systems, with a supplementary chapter by KATOK and L. MENDOZA, *Encyclopedia of Mathematics and its Applications* 54. Cambridge University Press, Cambridge, 1995.

- [Ki1] Y. KIFER, Large deviations in dynamical systems and stochastic processes, *Trans. Amer. Math. Soc.* 321:2 (1990), 505–524.
- [Ki2] Y. KIFER, Large deviations, averaging and periodic orbits of dynamical systems, *Comm. Math. Phys.* 162:1 (1994), 33–46.
- [L1] S.P. LALLEY, Renewal theorems in symbolic dynamics, with applications to geodesic flows, non-Euclidean tessellations and their fractal limits, *Acta Math.* 163:1-2 (1989), 1–55.
- [L2] S.P. LALLEY, Closed geodesics in homology classes on surfaces of variable negative curvature, *Duke Math. J.* 58:3 (1989), 795–821.
- [Le] G. LEBEAU, Équation des ondes amorties, *Algebraic and Geometric Methods in Mathematical Physics (Kaciveli, 1993)*, *Math. Phys. Stud.* 19, Kluwer Acad. Publ., Dordrecht (1996), 73–109.
- [Li] A.N. LIVŠIĆ, Certain properties of the homology of Y -systems, (in Russian) *Mat. Zametki* 10 (1971), 555–564; *Math. Notes* 10 (1971), 758–763.
- [MM] A.S. MARKUS, V.I. MATSAEV, Comparison theorems for spectra of linear operators and spectral asymptotics, *Trudy Moskov. Mat. Obshch.* 45 (1982), 133–181.
- [PP] W. PARRY, M. POLLICOTT, Zeta functions and the periodic orbit structure of hyperbolic dynamics, *Astérisque* 187/188 (1990).
- [PhR] R. PHILLIPS, Z. RUDNICK, *The circle problem in the hyperbolic plane*, *J. Funct. Anal.* 121:1 (1994), 78–116.
- [PoS] M. POLLICOTT, R. SHARP, Asymptotic expansions for closed orbits in homology classes, *Geom. Dedicata* 87:1-3 (2001), 123–160.
- [R] M. RATNER, The central limit theorem for geodesic flows on n -dimensional manifolds of negative curvature, *Israel J. Math.* 16 (1973), 181–197.
- [ReS] M. REED, B. SIMON, *Methods of Modern Mathematical Physics I. Functional Analysis*, Second edition. Academic Press, Inc. [Harcourt Brace Jovanovich, Publishers], New York, 1980.
- [Ru] W. RUDIN, *Real and Complex Analysis*, Third edition, McGraw-Hill Book Co., New York, 1987.
- [Rue] D. RUELLÉ, *Thermodynamic Formalism*, *Encyclopedia of Mathematics and its Applications* 5, Addison-Wesley Publishing Co., Reading, Mass., 1978.
- [S] E. SCHENCK, Energy decay for the damped wave equation under a pressure condition, [arXiv:0909.2093](https://arxiv.org/abs/0909.2093)
- [Sc] A. SCHNIRELMAN, Ergodic properties of eigenfunctions, *Usp. Math. Nauk.* 29 (1974), 181–182.
- [Se] A. SELBERG, Harmonic analysis and discontinuous groups in weakly symmetric Riemannian spaces with applications to Dirichlet series, *J. Indian Math. Soc.* 20 (1956), 47–87.
- [Sj] J. SJÖSTRAND, Asymptotic distribution of eigenfrequencies for damped wave equations, *Publ. Res. Inst. Math. Sci.* 36:5 (2000), 573–611.

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